

NAVAL POSTGRADUATE SCHOOL Monterey, California

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THESIS

"HEAD-ON" SCATTERING OF A TUBULAR CYLINDER OF FINITE LENGTH FOR RADAR TARGET IDENTIFICATION PURPOSES

by

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March 1985

Thesis Advisor:

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At this aspect angle, the back scattered fields depend only on the first Fourier component of the circumferential variations of the ϕ -current.

Measurements of several scaled tubular cylinders were taken and the experimental results were compared to theoretical data available.

This thesis is part of an ongoing project of target identification through the investigation of the cross section of a target over a broad frequency band. Approved for public release; distribution is unlimited.

"Head-on" Scattering of a Tubular Cylinder of Finite Length for Radar Target Identification Purposes.

bу

David Geller
Lieutenant, Israeli Navy
B.S., Technion-Israeli High Institute of Technology, 1977

Submitted in partial fulfillment of the requirements for the degree of

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ABSTRACT

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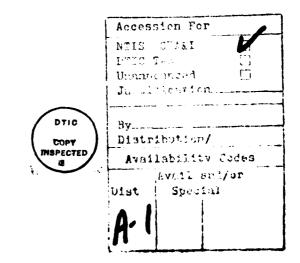


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I. INTRODUCTION

Identifying a target by its back scattering signal is desirable if different actions are to be taken towards different targets. At a time when warfare is no longer conducted face to face, but missiles have the ability to destroy thier targets long before they come within visual range, it has become necessary to identify targets by some means other than vision. The identification is needed because of the fact that current operational policy requires a positive identification of a target before destroying it. This in a sense completly negates the prime capability of sophisticated weapon-systems. For example, the HARPOON missile whose range is far beyond the horizon prevents a visual identification of the target from the fireing platform. A fact that limits the use of it to only specific situations under sever restrictions.

At present, most identifications are done with radars with which the presence of a target can be detected. In addition to the fact that the presence of a target can be discovered, information about its position, speed and acceleration may be obtained. An experienced radar operator can sometimes distinguish a big target from a small target by its spot-size on the radar screen, but this is not enough to identify the target.

To enhance our capability of identifying a target, it is advantageous to look at resonances exited by electromagnetic fields of different frequencies incident on the target. The reasons for looking at this particular range of frequencies are the following: First, at resonance the scattered fields are stronger compared to non-resonance situations. The

back-scattering cross-sections are larger and the target can be detected more easily. Second, the resonance frequencies and the amplitudes and phase shifts of the scattered fields at these frequencies are determined by the geometry of the target. Since the number of targets of our interest are finite, the identity of a target is revealed by examing a small data base.

This thesis is part of an ongoing project of target-identification through the investigation of the wide-band cross-section of a target. A finite tubular cylinder which has a circular cross section and a very thin conducting wall serves as the canonical target. The finite tubular cylinder is chosen because of its resemblance to a missile body and because theoretical formulations are available for its surface current distribution and scattered fields.

The tubular cylinder was placed in an anechoic chamber and was irradiated with an incident electromagnetic wave. While irradiated by the incident electromagnetic wave, surface currents were excited and generated scattered-field. The surface current had an axial and a circumferential component. The circumferential current circled around the cylinder while the axial current traveled along the cylinder and was reflected at the ends. At the same time, the incident wave kept impinging on the cylinder and excited new surface currents which added to the existing ones. certain frequency, the newly excited axial current might add constructively to the current reflected from one end, resulted in a large axial current, which gave a strong scattered field. Or the newly excited circumferential current might add constructively to the current which had made a complete circle around the cylinder, and a strong scattered field could be observed.

This thesis studied the "head-on" back scattering of the cylinder. At this aspect angle, the back scattered fields depended only on the first Fourier component of the circumferential variation of the ϕ -current. Measurements of several scaled tubular cylinders were taken and the experimental results were compared to theoretical data available.

For the next step of this project, more complicated models evolving from a tubular cylinder to a missile will be constructed and studied. Fins will be added, one end of the cylinder will be closed and rounded to perturb the model further and finally wings will be added to make a true missile model. By comparing the scattering data of the models to those of the tubular cylinders, the effects of the successive perturbations to the physical structure on the back-scattering cross-section and phase shift will be investigated.

Chapter II deals in its first part with electromagnetic back scattering theory in general, the definition of radar cross section, the polarization matrix and methods to obtain it. Its second part contains the solution to the electromagnetic back-scattering of a tubular cylinder with finite length. Chapter III describes some CW step frequency cross section measurements which are carried out, including the experimental setup, the measurement procedures and the measured results. Chapter IV deals with data analysis, compares the experimental results to theoretical data and presents some conclusions and recommendations for the future.

II. ELECTROMAGNETIC SCATTERING THEORY

A typical problem in electromagnetic scattering consists of these main elements.

- 1. Transmitting system: RF source and antennas.
- 2.An object of arbitrary shape and size as a target.
- 3. Receiving system: Antenna and receiving equipment to determine the amplitude, phase and polarization of fields at any point in space.

The transmitting system causes incident fields E^1 , H^1 to imping upon the target. The current in the source induces time-varying distributions of oscillating charges and currents in the scatterer. These currents cause scattered, or reradiated fields E^1 , H^1 . The total fields E^1 , H^1 are the vector sum of the incident and scattered fields.

In order to simplify the theoretical problem, it is usually assumed that the source is not coupled to the target. This fact enables one to obtain the scattered field by subtracting the incident fields from the total fields.

A configuration of the scattering problem is shown in Figure $2.1\,$

This chapter deals with the analytical background, the definition of radar cross-section, the polarization scattering matrix, and finally the specific problem of the scattering by a finite tubular cylinder.

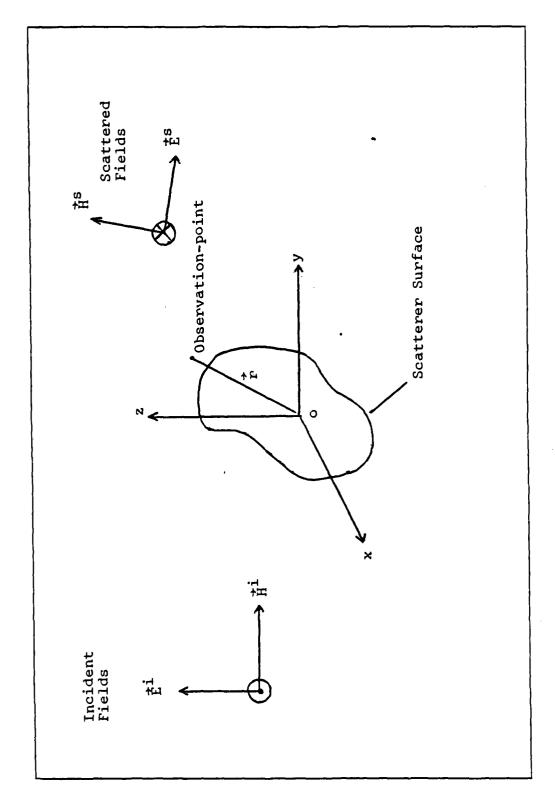


Figure 2.1 The Scattering Problem

A. ANALYTICAL BACKGROUND

1. Definition of Radar Cross-section

The radar cross-section of a target is a quantitative measure of the ratio of power density in the vector signal scattered in the direction of the receiver to the power density of the radar wave incident upon the target. The vectorial nature of the electromagnetic interaction requires specifications of the polarization of the incident wave with reference to target orientation in three dimensions. Radar operating frequency is an additional parameter which must be specified.

Thus, a single number specification for the radar cross section holds for a particular target, a specific polarization and the frequency of the incident wave, the aspect angle of the target relative to the incident wave, and the polarization of the receiving antenna.

The radar cross-section is defined to be independent of range to the target. This definition holds under far field assumption. Which means that the target is sufficiently far from the transmitting antenna to justify the assumption that the incident wave is planar at the target and the scattered wave is also planar in the neighborhood of the receiving antenna.

The theoretical definition of the radar crosssection relates incident to scattered electromagnetic fields, as given in equation 2.1

$$\sigma = 4\pi R^2 \lim_{R \to \infty} \left| \frac{E_s}{E_o} \right|^2$$
 (eqn 2.1)

where:

E_O = magnitude of electric field component of incident electromagnetic field at the target.

E_S = magnitude of electric field component of scattered electromagnetic field as measured by a hypothetical observer.

R = distance from target to the hypothetical observer.

The radar cross-section $\boldsymbol{\sigma}$ has the dimensions of area $\boldsymbol{m^2}$.

The limiting process is introduced in equation 2.1 to assure that the distance at which the hypothetical observation is made, is far enough from the target. Under the free-space conditions assumed, the quantity $|\mathbf{E}_{\mathrm{S}}/\mathbf{E}_{\mathrm{O}}|^2$ is proportional to the power flux density of the scattered waves.

2. Polarization Scattering Matrix

The radar cross-section of a target depends upon the target shape and material, the angle (or angles, in a case of bistatic system) at which the target is viewed, radar frequency, and the polarization of the radar-transmitting and receiving antennas.

In particular, if a target is viewed at a specific aspect angle with a single frequency, the radar cross section depends upon polarization [Ref. 1].

The polarization scattering matrix is introduced in order to express target reradiation independent of radar polarization [Ref. 2].

Scattering is expressed as an explicit function of radar polarization, when matrices describing the polarization properties of antennas and target are defined.

The transmitting and receiving antennas can be represented by the matrices:

$$\hat{q} = \begin{bmatrix} \cos \phi_t \\ \sin \phi_t e^{j \delta t} \end{bmatrix}$$
 (eqn 2.2)

$$\hat{p} = [\cos\phi_r \quad \sin\phi_r e^{j\delta r}] \qquad (eqn 2.3)$$

where:

q = column matrix defining the polarization of the transmitting antenna.

 \hat{p} = row matrix defining the polarization of the receiving antenna.

 ϕ = an angle which when δ =0, denotes the orientation of the linear polarization of the antenna when used for transmitting, refered to the horizontal plane.

 δ = phase angle.

t = denotes transmitting antenna.

r = denotes receiving antenna.

The configuration of the cylinder for the measurement in this thesis and the antenna polarization angle are shown in Figure 2.2

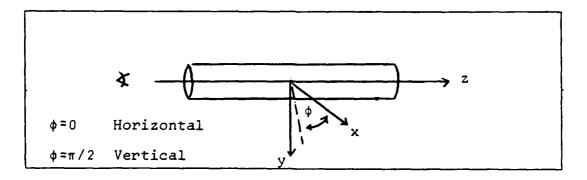


Figure 2.2 The Antenna Polarization Angle

The radar cross-section of a target observed by a transmitting antenna with polarization q and a receiving antenna with polarization p is given by equation 2.4

$$\sigma = |\hat{p}S\hat{q}|^2 \qquad (eqn 2.4)$$

where S denotes the complex scattering matrix used to represent the polarization properties of the target. Assuming a uniform plane incident wave, the scattering matrix is a linear relation between the incident field and the scattered far field from the target.

With p and q defind by equations 2.2 and 2.3 the scattering matrix S is a $2 \div 2$ matrix as shown below in equation 2.5

$$S = \begin{bmatrix} \sqrt{\sigma_{HH}} & 1^{j\rho_{HH}} & \sqrt{\sigma_{HV}} & 1^{j\rho_{HV}} \\ \sqrt{\sigma_{VH}} & 1^{j\rho_{VH}} & \sqrt{\sigma_{VP}} & 1^{j\rho_{VV}} \end{bmatrix}$$
 (eqn 2.5)

where:

 $\sqrt{\sigma}$ = magnitude of the scattering matrix element.

 ρ = phase of the scattering matrix element.

H = denotes horizontal polarization.

V = denotes vertical polarization.

The scattering matrix is symmetrical $(\sqrt{\sigma_{HV}} \sqrt{\sigma_{VH}}, \rho_{HV} = \rho_{VH})$ in at list two cases .

1. Bistatic scattering when the body is a perfect conductor.

2. Back-scattering from an arbitrary body.

The transmitting and receiving antennas used in our case are linearly horizontally polarized.

The expressions for the transmitting and receiving antennas under this condition are:

$$\hat{q} = \begin{bmatrix} 1 \\ 0 \end{bmatrix}$$
 (eqn 2.6)

$$\hat{p} = [1 \ 0]$$
 (eqn 2.7)

since $\phi = 0$.

The radar cross-section is:

$$\sigma = \sigma_{HH}$$
 (eqn 2.8)

3. Methods of Obtaining the Scattering Matrix

Ideally one would compute the radar cross-section of a target through the formal solution of Maxwells equations. Those equations should be solved for the boundary conditions appropriate to the target.

The major mathematical method for obtaining an exact solution is the separation of variables.

Formal solutions via separation of variables are possible only for a few special cases. In those cases the wave equation is separable in a coordinate system, having a

coordinate surface that coincides with the surface of the body [Ref. 3]. This situation explains why exact solutions are rare, and for many practical problems, the use of approximations is the only practical approach.

The integral equation formulation shows that electromagnetic scattering of an incident plane wave by an arbitrary body can be described in terms of an integral of various vector products. Those vector products involve the surface electric and magnetic fields. One form which is convenient for this purpose is the Chu-Stratton integral. [Ref. 4]. This integral is an exact representation of the scattered electromagnetic field. It is given in terms of an integration over a complete surface enclosing the body in question. In particular, if there were available knowledge of the total distribution of electric and magnetic fields about the body, insertion of these values in the Chu-Straton integral would permit the immediate solution of the scattering problem.

Numerical schemes have been designed to solve the integral equations approximately. High speed digital computers are needed to establish surface currents flowing on the target. For example surface current distribution can be computed by a finite difference solution to a network of simultaneous equations [Ref. 5].

Because of limitation of computation time and storage capability, the finite difference solution is applicable only when the dimensions of the target do not exceed a very few wavelengths.

For targets larger than a few wavelengths in dimension, asymptotic methods are frequently used. There are three levels of complexity:

The simplest approach is the geometric optics approach. It treats ray bundles by the laws of reflection and refraction [Ref. 6]. However, geometric-optics failes

to distinguish the effects of polarization and the wave nature of the problem.

The second approach is the physical optics. In this approach the local current density, at each point on the illuminated portion of the body, is assumed to be identical to that which would flow at that point on an infinite tangent plane [Ref. 7]. Physical optics is not valid for applications entailing accurate specification effects [Ref. 8].

The third approach is the geometric diffraction theory. This approach is an extension of geometric optic that accounts for diffraction [Ref. 9]. The technique combines the simplicity inherent in the ray approach with the necessary consideration of wavelengths and phases of the wave. But it never include resonances of the target.

B. SCATTERING BY A FINITE TUBULAR CYLINDER

When a cylinder is irradiated with an incident electromagnetic wave, surface current is exited on the cylinder. This surface current will radiate and generate the scattered field.

Electromagnetic scattering from conducting objects in a homogeneous, isotropic medium can be treated as a boundary value problem. By use of the Stratton-Chu equations, integrodifferential equations can be set up for the current distribution on the surfaces of the objects, with the Greens function in the medium as the kernel.

The current distribution on the surface of a tubular cylindrical conductor with negligible wall thickness, excited by an incident electromagnetic field, can be written as a pair of coupled integrodifferential equations with the sum of the inside and outside surface currents as the unknown and the incident tangential electric field on

the surface of the conductor as the given quantity. Such a coupled integrodifferential equations were given by Lee [Ref. 16].

$$(1 + \frac{1}{1_1^2} \frac{3^2}{3^2}) \int_{-1}^{1} dz_0 K_{z,n}(z_0) G_n(1_1 | z - z_0 | , 1_2) \quad (eqn \ 2.9)$$

$$+ \frac{\text{in}}{l_1 l_2} \frac{\partial}{\partial z} \int_{-1}^{1} dz_0 K_{\phi n}(z_0) G_n(l_1 | z - z_0 |, l_2) = - \frac{2i}{l_1 l_2 \xi_0} E_{zn}^{sc}(z).$$

$$\int_{-1}^{1} dz_{o} K_{\phi n}(z_{o}) \{ \frac{1}{2} [G_{n-1}(1_{1}|z-z_{o}|,1_{2})$$
 (eqn 2.10)

+
$$G_{n+1}(1_1|z-z_0|,1_2)] - \frac{n^2}{1_2^2} G_n(1_1|z-z_0|,1_2)$$

$$+ \frac{in}{l_1 l_2} \frac{\partial}{\partial z} \int_{-1}^{1} dz_{o} K_{zn}(z_{o}) G_{n}(l_1 | z - z_{o}|, l_2) = - \frac{2i}{l_1 l_2 \xi_{o}} E_{\phi n}(z).$$

equations 2.9 and 2.10 can be obtained from the Stratton-Chu equations, together with the edge conditions that:

$$K_z(0,z) = 0 \cdot (1-z^2)^{1/2}$$
 as $|z| \to 1$ (eqn 2.11)

In the previous equations:

sc-denotes-scattered.

i -denotes-incident.

 $\xi = (\mu/\epsilon)^{1/2}$

and

$$G_{n}(l_{1}|z-z_{0}|, l_{2}) = (eqn 2.12)$$

$$\int_{-\pi}^{\pi} \frac{d\phi}{2\pi} l^{-in(\phi-\phi_{0})} \cdot G[l_{1}|z-z_{0}|, 2l_{2}|\sin(\phi-\phi_{0})/2|]$$

$$G(x_{1},x_{2}) = \{exp[i(x_{1}^{2}+x_{2}^{2})^{1/2}]\}(x_{1}^{2}+x_{2}^{2})^{1/2} (eqn 2.13)$$

The following equations, 2.14 and 2.15 are boundary conditions for the tangential electric field components on a perfectly conducting surface.

$$E_{zn}^{SC}(z) + E_{zn}^{i}(z) = 0$$
 -1

$$E_{\phi n}^{SC}(z) + E_{\phi n}^{i}(z) = 0$$
 -1

The cylinder is assumed to be in a medium with homogeneous isotropic permittivity ε and permeability μ . The half length of the cylinder l_1 and the radius l_2 are measured in 1/k, with $k=\omega\left(\varepsilon\mu\right)$ being the wave number, so that $l_1=kh$ and l_7ka .

The coordinate system is scaled so that the tubular cylinder occupies the region -1 < z < 1, $\rho = 1$ as shown in Figure 2.3

On the surface, ρ =1,there are scattered electric field $E^{SC}(\phi,Z)$, the incident electric field $E^{\dot{i}}(\phi,Z)$ and for -1<Z<1,the surface current $\ddot{K}(\phi,Z)$,which is the sum of the

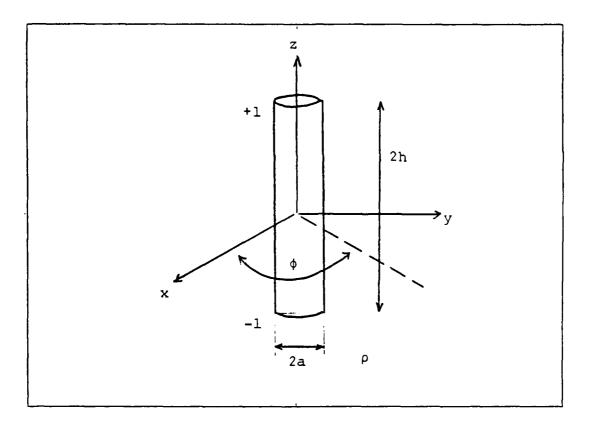


Figure 2.3 The Tubular Cylinder

outer surface current $\vec{K}^{\dagger}(\phi,Z)$ on $\rho=1^{\dagger}$ and the inner surface current $\vec{K}^{\dagger}(\phi,Z)$ on $\rho=1^{\dagger}$.

The surface currents can be represented as:

(eqn 2.16)

$$K_{z}(\phi,z) = \sum_{n=-\infty}^{\infty} e^{jn\phi} K_{zn}(z) = \sum_{n=0}^{\infty} [K_{zn}^{(+)}(z) \cosh \phi + iK_{zn}^{(-)}(z) \sinh \phi].$$

where

$$K_{zo}^{(+)}(z) = K_{zo}(z).$$

$$K_{Z_{n}}^{(+)}(z) = K_{Z_{n}}(z) + K_{Z_{n}-n}(z)$$
 $K_{Z_{n}}^{(-)}(z) = 0$

 $K_{\phi}(\phi,z) = \sum_{n=-\infty}^{\infty} e^{jn\phi} K\phi_{n}(z) = \sum_{n=0}^{\infty} \left[K_{\phi}^{(+)}(z) \cosh\phi + iK_{\phi}^{(-)}(z) \sinh\phi\right].$

where

$$K_{\phi,0}^{(+)}(z) = K_{\phi,0}(z).$$

$$K_{\phi}^{(+)}(z) = K_{\phi,n}(z) + K_{\phi,-n}(z)$$
.

$$K_{\phi,0}^{(-)}(z) = 0.$$

 $K_Z^{}(\varphi,z)$ is the axial current density,and $K_\phi^{}(\varphi,z)$ is the circumferential current density.

In the far field where:

$$K|\dot{r}| = (1_2^2 \rho^2 + 1_1^2 z^2)^{1/2} >> 1$$
 (eqn 2.18)

and

$$|K|\hat{r}| > (1_1^2 + 1_2^2)^{1/2} \ge |K|\hat{r}_0|$$
 (eqn 2.19)

$$-\frac{2i}{l_{1}l_{2}\xi_{o}} E_{z}(\rho,\phi,z) = \sin^{2}\theta \int_{-1}^{1} dz_{o} \int_{-\pi}^{\pi} \frac{d\phi_{o}}{2\pi} G(\vec{r}-\vec{r}_{o}) K_{z}(\phi_{o},z_{o}) - \sin\theta\cos\theta \int_{-1}^{1} dz_{o} \int_{-\pi}^{\pi} \frac{d\phi_{o}}{2\pi} G(\vec{r}-\vec{r}_{o}) \sin(\phi-\phi_{o})$$
 (eqn 2.20)
$$K_{\phi}(\phi_{o},z_{o})$$

$$\frac{-2i}{l_1 l_2 l_0^2} \operatorname{Ep}(\rho, \phi, z) = -\sin\theta \cos\theta \int_{-1}^{1} dz_0 \int_{-\pi}^{\pi} \frac{d\phi_0}{2\pi} G(r - r_0) Kz(\phi_0, z_0) + \cos^2\theta \int_{-1}^{1} dz_0 \int_{-\pi}^{\pi} \frac{d\phi_0}{2\pi} G(r - r_0) \sin(\phi - \phi_0) K_{\phi}(\phi_0, z_0).$$
(eqn 2.21)

$$\frac{-2i}{l_1 l_2 \xi_0} E_{\phi}(\rho, \phi, z) = \int_{-1}^{1} dz_0 \int_{-\pi}^{\pi} \frac{d\phi_0}{2\pi} G(\vec{r} - \vec{r}_0) \cos(\phi - \phi_0) K_{\phi}(\phi_0, z_0)$$

(eqn 2.22)

In spherical coordinates, the expressions are simpler. Because:

$$E_{r}(r,\theta,\phi)=E_{0}(\rho,\phi,z)\sin\theta+Ez(\rho,\phi,z)\cos\theta$$
 (eqn 2.23)

$$E_{\theta}(r,\theta,\phi)=E_{\rho}(\rho,\phi,z)\cos\theta-E_{z}(\rho,\phi,z)\sin\theta$$
 (eqn 2.24)

equations 2.20, 2.21 and 2.22 can be written in the following form:

$$-\frac{2i}{l_1 l_2 \xi_0} E_{\theta}(r, \theta, \phi) = 0$$
 (eqn 2.25)

$$-\frac{2i}{1_{1}^{1}_{2}\xi_{o}} E_{\theta}(r,\theta,\phi) = -\sin\theta \int_{-1}^{1} dz_{o} \int_{-\pi}^{\pi} \frac{d\phi_{o}}{2\pi} G(r-r_{o})K_{z}(\phi_{o},z_{o}) + (eqn 2.26)$$

$$\cos\theta \int_{-1}^{1} dz_{o} \int_{-\pi}^{\pi} \frac{d\phi_{o}}{2\pi} G(\vec{r} - \vec{r}_{o}) \sin(\phi - \phi_{o}) K_{\phi}(\phi_{o}, z_{o})$$

$$-\frac{2i}{l_1 l_2 \xi_0} E_{\phi}(r,\theta,\phi) = \int_{-1}^{1} dz_e \int_{-\pi}^{\pi} \frac{d\phi_e}{2\pi} G(r-r_0) \cos(\phi-\phi_0) K_{\phi}(\phi_0,z_0)$$

(eqn 2.27)

where

$$G(\vec{r}-\vec{r}_{o}) = G(\vec{r})\exp^{-il_{2}\sin\theta\cos(\phi-\phi_{o})}\exp^{-il_{1}\cos\theta z_{o}}$$

since:

$$\int_{-\pi}^{\pi} \frac{d\phi}{2\pi} \cosh \exp \frac{-il_2 \sin\theta \cos(\phi - \phi_0)}{=i^{-n} Jn(l_2 \sin\theta)}$$

(eqn 2.30)

$$-1 \int_{-1}^{1} \frac{dz_{e}}{\pi \sqrt{1-z_{o}^{2}}} cospve^{-il_{1}cosvz} = \int_{0}^{\pi} \frac{dv}{\pi} cospve^{-il_{1}cos\theta} cosv$$
$$= i^{-p} Jp(l_{1}cos\theta)$$

and with z=cosv and the facts that $k_{\phi} \rightarrow (1-Z^2)^{4/2}, K_z \rightarrow (1-z^2)^{4/2}$ on the edges of the cylinder

$$K_{zn}(z) = \frac{1}{\pi} \sum_{p=0}^{\infty} K_{z,n}^{p} \sin(p+1)v$$
 (eqn 2.31)

$$K_{\phi n}(z) = \frac{1}{\pi \sin v} \sum_{p=0}^{\infty} K_{\phi,n}^{p} \cos pv \qquad (eqn 2.32)$$

where $K_{z\,,\,n}^{\,p}$ and $K_{\varphi\,,\,n}^{\,p}$ the sum of the inside and outside currents from equation 2.16

$$K_{z,n}^{(+)p} = K_{z,n} + K_{z,-n}$$
 (eqn 2.33)

$$K_{\phi}^{(+)p} = K_{\phi,n} + K_{\phi,-n}$$
 (eqn 2.34)

The final equations for the general case are:

$$E_{r}(r,\theta,\phi) = 0 \qquad (eqn 2.35)$$

$$-\frac{2i}{l_1 l_2 \xi_0 G(r)} E_{\theta}(r,\theta,\phi) = -\sum_{n=0}^{\infty} \sum_{p=0}^{\infty} i^{-(n+p)} \frac{(p+1)\sin\theta}{l_1 \cos\theta}$$
(eqn 2.36)

$$J_{p+1}(l_1\cos\theta)J_n(l_2\sin\theta)[K_{zn}^{(+)p}\cos n\phi+iK_{zn}^{(-)p}\sin n\phi]$$

$$-\frac{2i}{l_1 l_2 \xi G(r)} E_{\phi}(r,\theta,\phi) =$$

(eqn 2.37)

Since $E_r(r,\theta,\phi)=0$ in the far field

(eqn 2.38)

$$E_{x}(r,\theta,\phi)=E_{\theta}(r,\theta,\phi)\cos\theta\cos\phi-E_{\theta}(r,\theta,\phi)\sin\phi$$

(eqn 2.39)

$$E_{y}(r,\theta,\phi) = E_{\theta}(r,\theta,\phi)\cos\theta\sin\phi + E_{\phi}(r,\theta,\phi)\cos\phi$$

(eqn 2.40)

$$E_z(r,\theta,\phi) = -E_{\theta}(r,\theta,\phi)\sin\theta$$

In the special case, with the cylinder positioned head-on to the antennas, the incident and scattered fields can be described as follows:

 $\theta = \pi$ (along the Z axis)

The scattered fields are:

(eqn 2.41)

$$\frac{-2i}{l_{1}l_{2}\xi_{0}G(r)}E_{\theta}(r,\pi,\phi) = \frac{1}{2}\sum_{p=0}^{\infty} i^{-(p+1)}J_{p}(l_{1})[K_{\phi 1}^{(-)p}\cos\phi + iK_{\phi 1}^{(+)p}\sin\phi]$$

(eqn 2.42)

$$\frac{-2i}{1_{1}1_{2}\xi_{0}G(\dot{r})}E_{\phi}(r,\pi,\phi)=\frac{1}{2}\sum_{p=0}^{\infty}i^{-(p+1)}J_{p}(1_{1})[-K_{\phi 1}^{(-)p}sin\phi+iK_{\phi 1}^{(+)p}cos\phi].$$

or in rectangular coordinates, along the -z axis:

(eqn 2.43)

$$\frac{-2i}{l_1 l_2 \xi_0 G(r)} E_{x}(z) = \frac{1}{2} \sum_{p=0}^{\infty} i^{(p-1)} J_p(l_1) K_{\phi 1}^{(-)p}$$

(eqn 2.44)

$$\frac{-2i}{l_1 l_2 \xi_0 G(r)} E_y(z) = \frac{1}{2} \sum_{p=0}^{\infty} i^{-p} J_p(l_1) K_{\phi 1}^{(+)p}$$

(eqn 2.45)

$$E_z(z) = 0.$$

The configuration of the cylinder and the incident field is shown in Figure 2.4

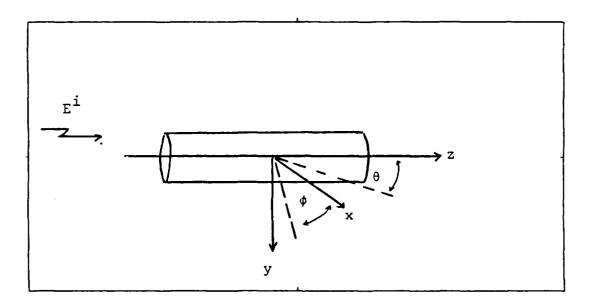


Figure 2.4 The Cylinder and the Incident Field

With the axis of the cylinder chosen along the z-direction and the incident field along the z axis, the incident field is given as:

$$\vec{E}^{inc} = \hat{e}_2 1^{ikz}$$
 , $\hat{e}_2 = -\hat{x}$ (eqn 2.46)

$$E_{\phi}^{\text{inc}} = +\sin\phi e^{ikz} \qquad (eqn 2.47)$$

Because E_{ϕ}^{inc} is an odd function in ϕ

$$K_{\phi n}^{(+)}(z) = 0$$

From equations 2.43- 2.45 $E_y(Z)=0$, and $\vec{E}(Z)=\hat{x}E_x(Z)$ along the -Z axis.

Cross-section and phase of a scatterer are defined with a linearly polarized plane incident wave on the scatterer. The incident wave has unit strength and zero phase at the center of the scatterer.

The cross-section is given by:

$$\sigma = \lim_{n \to \infty} 4\pi r^2 \qquad \left| E^{\text{SC}} \right|^2 \qquad (\text{eqn 2.48})$$

and the phase shift is given by:

$$\delta = \arg(1^{-ikr}E^{sc}) = \arg\left[\sum_{p=0}^{\infty} i^p J_p(1_1) K_{\phi 1}^{(-)p}\right]. \quad (eqn 2.49)$$

since:

$$|E^{SC}| = |\frac{1_1 1_2 \xi_0 G(r)}{4i} \sum_{p=0}^{\infty} i^{p-1} J_p(1_1) K_{\phi 1}^{(-)p}|$$
 (eqn 2.50)

the cross-section of the finite cylinder is:

$$\sigma = \lim_{r \to \infty} \frac{4\pi r^2}{r^2} \left| \frac{1_1 1_2 \xi_0 G(r)}{41} \sum_{p=0}^{\infty} i^{p-1} J_p(1_1) K_{\phi 1}^{(-)p} \right|^2$$

$$= \pi a^{2} \left| \frac{kh\xi}{a} \sum_{p=0}^{\infty} i^{p-1} J_{p}(1) K_{\phi 1}^{(-)p} \right|^{2}$$

III. MEASUREMENTS AND RESULTS

Radar cross-section estimation is as much art as science. The air of mystery that surrounds it will only be removed by an increase in understanding of how a target scatters energy incident upon it.

A complete knowledge of the scattering behavior is available only for few bodies. For these bodies we do not have exact solutions for their radar cross section, the best we can do is to provide approximate values. Such values should be checked against experimental measurements.

This chapter describes some CW step frequency crosssection measurements carried out at the Naval Postgraduate School, including the experimental setup, the measurement procedure and the measurements results.

Computer programs for calibration of the experimental setup and target measurements process are given in Appendix A. An explanation to the programs can be found in previous thesis from the Naval Postgraduate School [Ref. 15].

A. LABORATORY DESCRIPTION

1. System Configuration

The physical setup of the laboratory can be divided into five parts:

- -The RF source.
- -Transmitting and receiving antennas.
- -The anechoic chamber.
- -Target mount, targets.
- -Control and data processing equipment.

This setup is called a Radar range geometry and is designed for cross section measurements of models. The configuration of the entire system is shown in Figure 3.1.

A variety of radar range geometries have been developed during the years. Those radar range geometries were characterized by the way they have been designed to eliminate unwanted signals reflected from the foreground and background [Ref. 10]. An important component of the setup at the NavalPostgraduate School is the anechoic chamber [Ref. 11]. The anechoic chamber is used to approximate free space conditions in a closed environment. The anechoic chamber is enclosed with aluminium plates and internally lined with a radio frequency absorbing material. absorbing material provides the necessary attenuation to the reflections from the walls, floor and ceiling, and the aluminium surface provides protection against external sources of noise such as atmospheric noise, man made noise (radio, television, radar etc) and weather conditions. characteristics of the absorber material are specified in Appendix B.

The target is supported by a styrofoam stand. The reason for choosing the material, is to minimize coupling between the stand and the target. [Ref. 12].

The radar cross section range utilizes two identical horn antennas to approximate a back scattering system. Both antennas are horizontally polarized. The antennas are mounted on a removable panel located in the front wall of the anechoic chamber. The antennas can be adjusted in all three axis, providing beam steering towards the target. The basic antennas characteristics are given in Table I, and the configuration of the anechoic chamber is shown in Figure 3.2

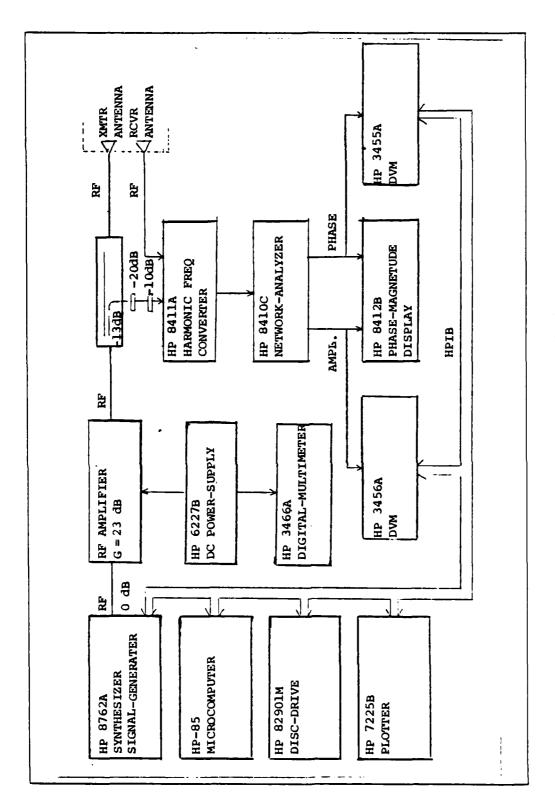


Figure 3.1 Block Diagram of the System

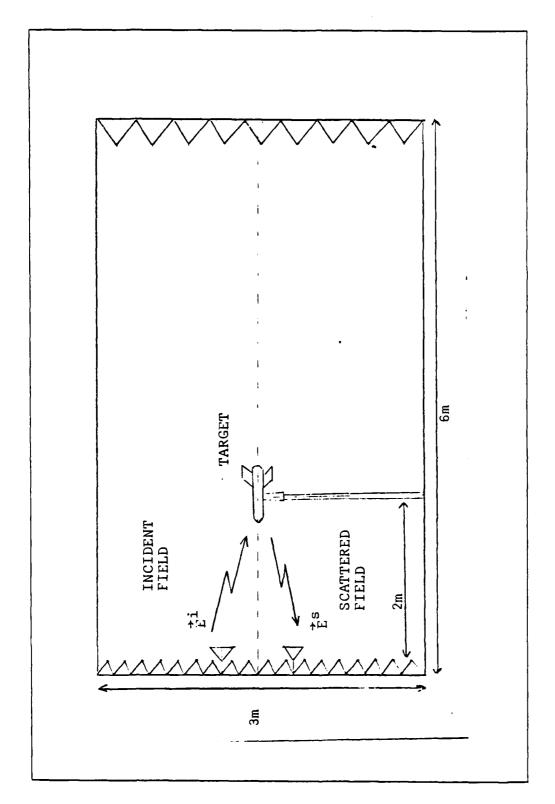


Figure 3.2 Configuration of the Anechoic Chamber

TABLE I ANTENNAS SPECIFICATION

Frequency range 4-18 GHz
Gain 6-12dB
VSWR (max) 3.2:1

Isolation between cross polarization <20dB below 5.5 GHz

>20dB above 5.5 GHz

2. Instrumentation

 \cdot The equipment of the experimental setup is listed in Table II .

The RF signal to the system is provided by the signal generator (HP-8672A). The RF output from the signal generator is amplified by an RF amplifier (Avantek-SA-83-2953). (The DC-power supply,HP-6227B,to the amplifier is monitored by a digital multimeter,HP-3466A).

The amplified RF signal passes the directional coupler (Narda-5292) and goes to the transmitting antenna. The directional coupler has a coupling coefficient of -13 dB and provides the reference signal to the harmonic frequency converter (HP-8411A) after the signal is further attenuated by 30 dB. The scattered signal from the target goes through the receiving antenna to the first port of the harmonic frequency converter. The frequency converter and the network analyzer (HP-8410C) convert the test channel signal into two 278 KHz signals containing the magnitude and phase information of the test channel signal relative to the refence channel signal. Both signals enter the phase-magnitude

TABLE II
LIST OF EXPERIMENTAL DEVICES

-Microcomputer	HP-85
-Synthesized Signal Generator	HP-8672A
-Harmonic Frequency converter	HP-8411A
-Network Analizer	HP-8410C
-Phase-Magnitude Display	HP-8412B
-RF Amplifier	Avantek sa-83-2953
-DC Power supply	HP-6277B
-Digital Multimeter	HP-3466A
-Digital voltmeter	HP-3455A
-Digital voltmeter	HP-3456A
-Flexible Disc Drive	HP-82901M
-Plotter	HP-7225B
-Directional Coupler	Narda 5292

display (HP-8412B) The DC plotter outputs of the display unit are fed to the two DVM's (HP-3455A, HP-3456a).

The complete setup is controlled and the data is processed by the microcomputer (HP-85) and the results are stored on discs.

3. Targets

The most direct means of obtaining knowledges about radar cross sections are by measurement of the radar return from the target itself or from an accurate model of the target [Ref. 13]. One advantage of a radar range is the practicability of testing models that are smaller and cheaper than full scale targets. The radar wavelength is scaled by the same factor as are the dimensions of the model. If D is any given dimension of the target, and D is the equivalent dimension of the model, the following scaling relation is employed:

$$D_{M}/D_{O} = \lambda_{M}/\lambda_{O}$$
 (eqn 3.1)

where λ_M is the wavelength used for the measurement and λ_0 is the wavelength for which the target radar cross-section is required.

At the same time, the measured cross section is altered in proportion to the change in power captured as a result of dimensional changes. If σ_M is the observed cross section of the model, the target cross section σ_0 is given by:

$$\sigma_{O} = \frac{\lambda_{O}^{2}}{\lambda_{M}^{2}} \sigma_{M} \qquad (eqn 3.2)$$

Exact scaling requires the model conductivity to be equal to the target conductivity multiplied by the ratio $(\lambda_0/\lambda_{\rm M})$, and model permittivity and permeability at the test frequency to equal corresponding target electrical properties at the operational frequency. It is necessary to make sure that the surface conductance of the model won't be smaller than the target.

The targets which are tested in the measurements are thin walled tubular cylinders made of brass. There are 20 cylinders of various lengths and diameters. The dimensions of the targets are given in Table III .

For the calibration of the system a 3.187 inch aluminium sphere is used.

Some measurements are taken with a cylinder with fins attached . The description of the cylinder with fins is shown in Figure 3.3 .

TABLE III
TARGET-CHARACTERISTICS

TARGET		LENGTH	DIAMETER	THICKNESS
cylinder	1	2.00"	0.375"	0.012"
cylinder	2	2.00"	0.500"	0.014"
cylinder	3	2.00"	0.750"	0.011"
cylinder	4	2.25"	0.375"	0.012"
cylinder	5	2.25"	0.500"	0.014"
cylinder	6	2.25"	0.750"	0.011"
cylinder	7	2.50"	0.375"	0.012"
cylinder	8	2.50"	0.500"	0.014"
cylinder	9	2.50"	0.750"	0.011"
cylinder	10	2.75"	0.375"	0.012"
cylinder	11	2.75"	0.500"	0.014"
cylinder	12	2.75"	0.750"	0.011"
cylinder	13	3.00"	0.375"	0.012"
cylinder	14	3.00"	0.500"	0.014"
cylinder	15	3.00"	0.750"	0.011"
cylinder	16	1.50"	0.375"	0.012"
cylinder	17	4.50"	0.750"	0.011"
cylinder	18	2.50"	0.625"	0.014"
cylinder	19	3.75"	0.750"	0.011"
cylinder	with fins	3.00"	0.750"	0.011"

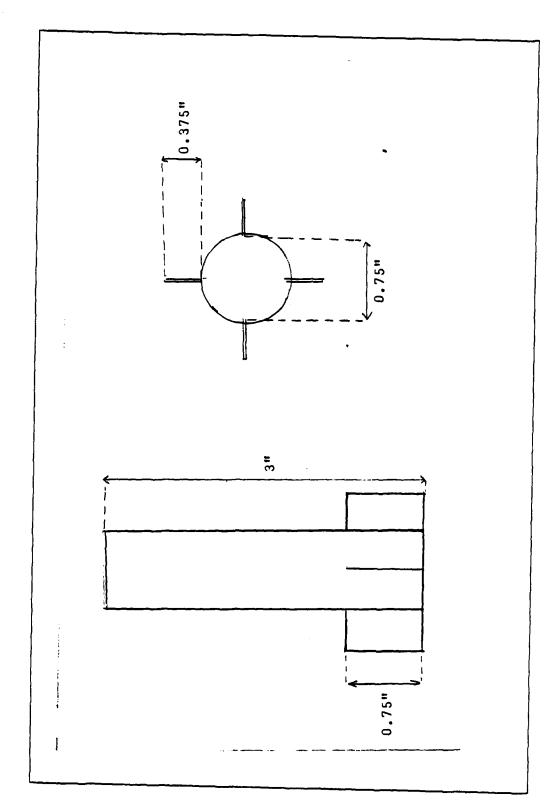


Figure 3.3 Tubular Cylinder with Fins

B. MEASUREMENT PROCEDURE

1. Calibration of the System

As a first step in the measurement, a calibration of the system must be done.

To calibrate the system output, a target of known cross section (usually a metal sphere) is placed at at the target support to fix the level of the calibration curve. This measurement assures that the entire system is calibrated in the proper frequency range.

Measurements are taken at many fixed frequencies between 10 to 15 GHz. Since both the transmitting and receiving antennas are horizontally polarized, the measurements are in one dimensional plane.

The calibration of the system is divided into two steps:

- 1. Take measurements without target in the anechoic chamber. (background data).
- 2. Inserte a 3.187 inch diameter aluminium sphere and repeat the measurements.

With the target in place, the received signal is a vectorial sum of the target echo and the background radiation. By taking background data, direct back scattering from the targets support and the walls of the anechoic chamber can be substracted from the vector sum.

The quality of the calibration of the system is tested and checked by comparing a new set of measured data on the sphere to thier theoretical values whenever a calibration process is done.

After the calibration of the system is accomplished, measurements of target cross sections can be started.

2. Measurements of the Targets

The measurements of the cylinders are taken in the frequency range of 10 to 15 GHz in steps of 0.1 GHz. The decision to limit the frequency range inspite of the ability of the equipment to operate beyond this range is due to the following reason:

Only in this frequency range consistent data can be obtained through averaging the data obtained from several frequency scans.

All the measurements are taken while the cylinders are positioned "Head-on" to the antennas as seen in Figure 3.4 A description of the cylinders tested is given in Table III, and reasons for choosing those particular dimensions for the cylinders are given in Chapter .

As a rule of thumb, far field approximations are good under the following conditions:

$$r > 10\lambda$$
 (eqn 3.3)

$$r > 10D$$
 (eqn 3.4)

$$r > \frac{2D^2}{\lambda}$$
 (eqn 3.5)

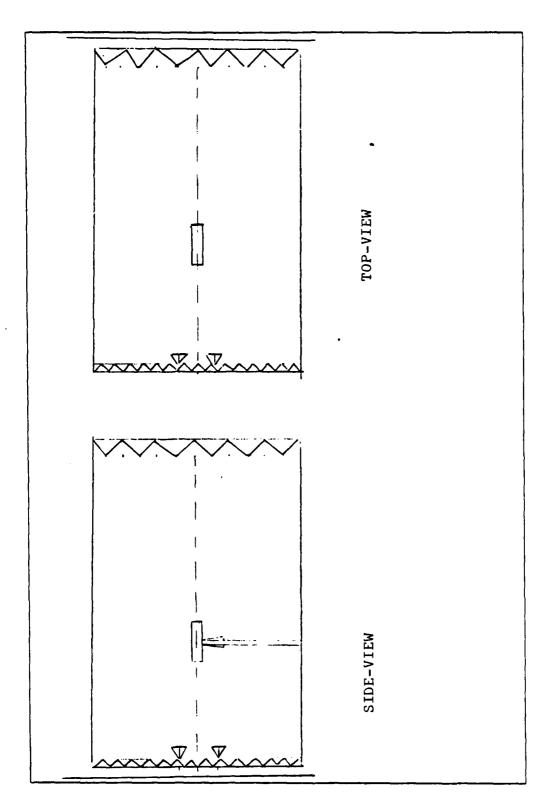


Figure 3.4 The Geometry Range

where r is the distance between the target and the antennas (receiving and transmitting), D is the largest dimension of either the target or the antennas and their separations, and λ is the wavelength. In our case, all conditions are met.

Before each target measurement, calibration of the system is done.

The experimental results for the targets in term of plots of their cross section and phases versus frequency are given at the end of the chapter.

3. Sources of Measurement Errors

Recognition of error sources in the measured data is difficult. Only in few special cases, it is possible to recognize the presence of an error, and to determine its source by observing the deviation of a target cross section versus frequency plot from anticipated behavior.

The errors in the measurement in our case are due to system noise and background noise.

The system noise, is caused mainly by the receiver. The network analyzer (HP-8410C) is a harmonic mixing receiver. It selects a harmonic of its internal VCO for the local oscillator frequency used to down convert the test frequency to the first IF. Harmonic skip errors can occur when the receiver selects a different harmonic (and VCO frequency) for the same frequency between system calibration and target measurement. The local oscillator power varies from one harmonic to another, thus the mixer transfer-characteristic varies. This fact causes random magnitude and phase variations of from 0.1 to 0.4 dB and up to 2 degrees.

As for the background noise, it is caused by coupling between the target and its mechanical support, and by strong coupling between the antennas in low frequencies.

In our case, the coupling between the antennas is reduced by going to higher range of frequencies. The

coupling between the target and its mechanical support is unavoidable and finally the receiver noise is reduced by averaging the measured results.

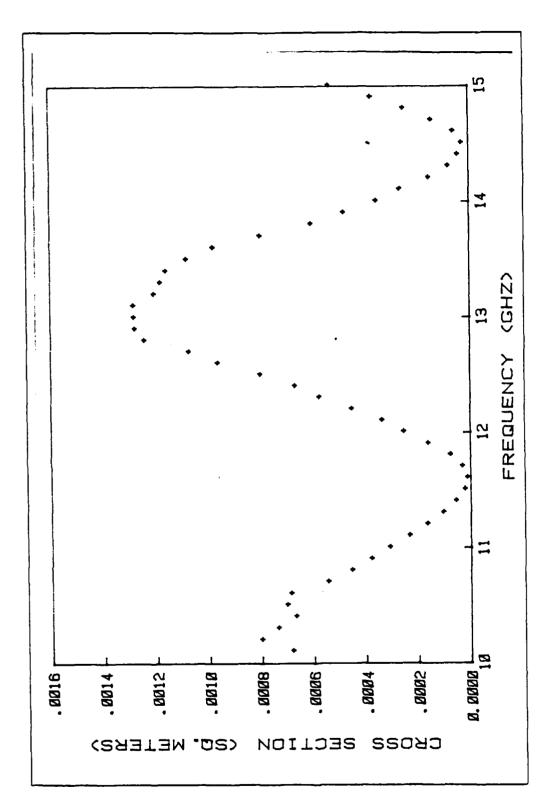


Figure 3.5 Cylinder 1: Cross-section

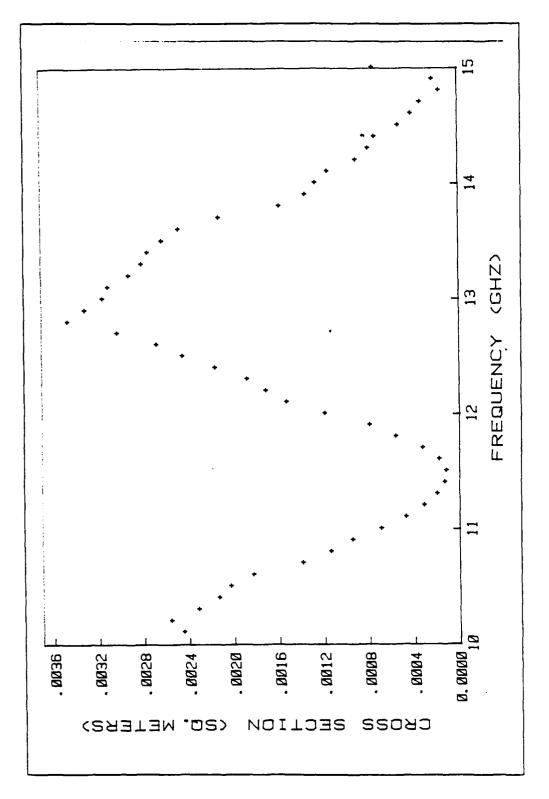


Figure 3.6 Cylinder 2: Cross-section

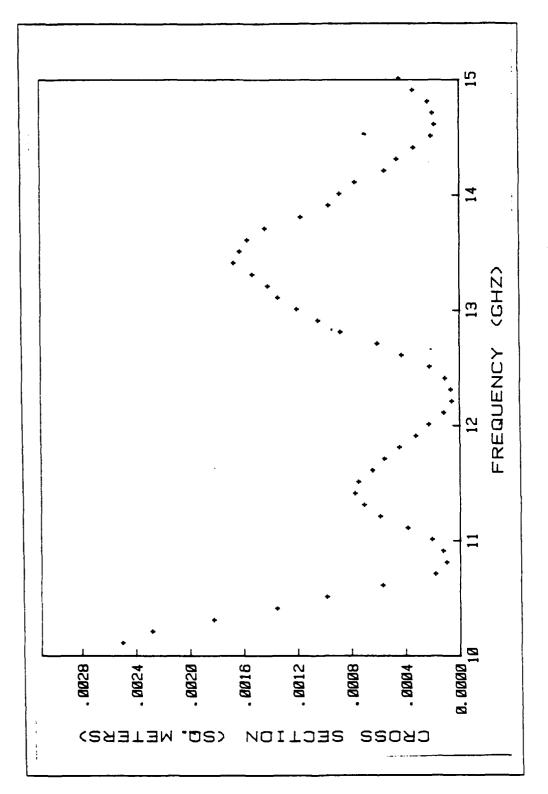


Figure 3.7 Cylinder 3: Cross-section

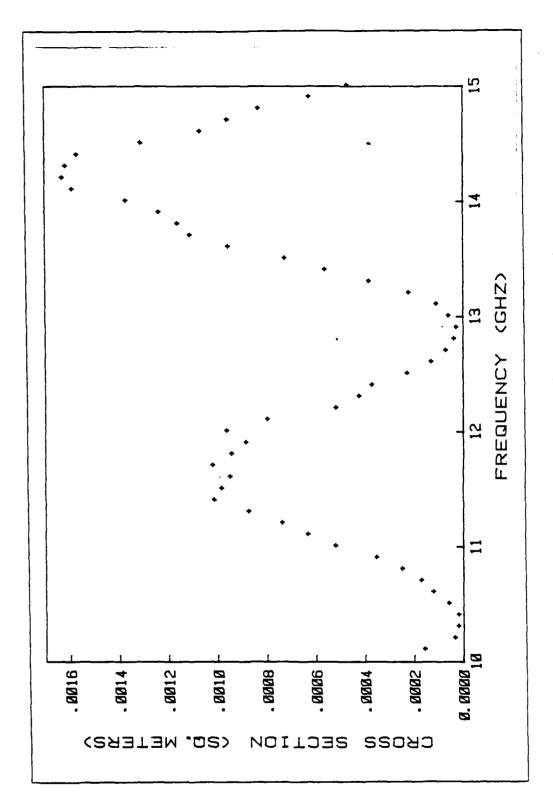


Figure 3.8 Cylinder 4: Cross-section

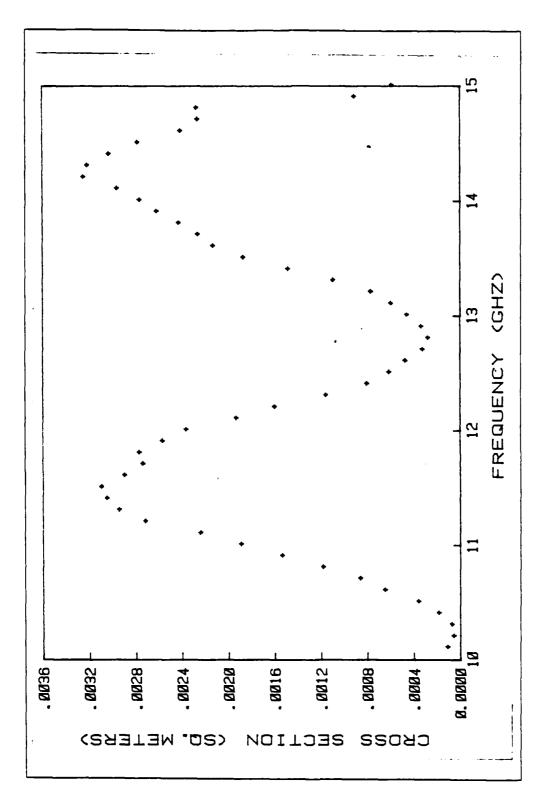


Figure 3.9 Cylinder 5: Cross-section

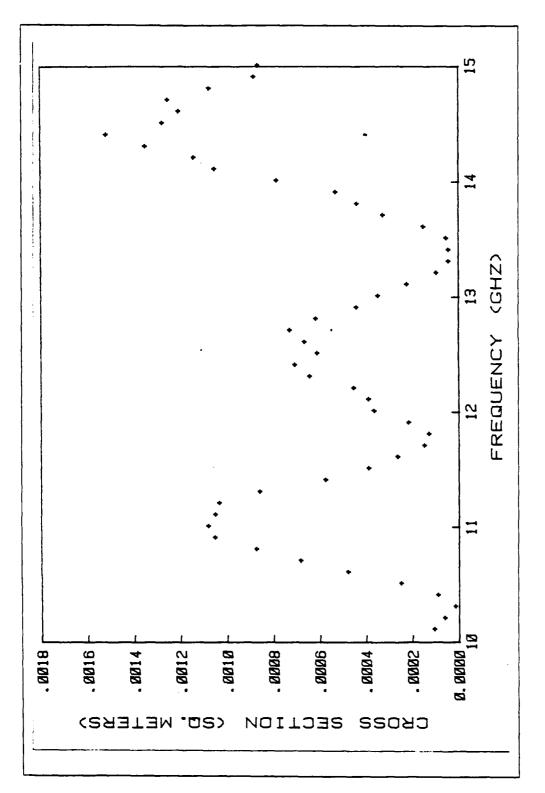


Figure 3.10 Cylinder 6: Cross-section

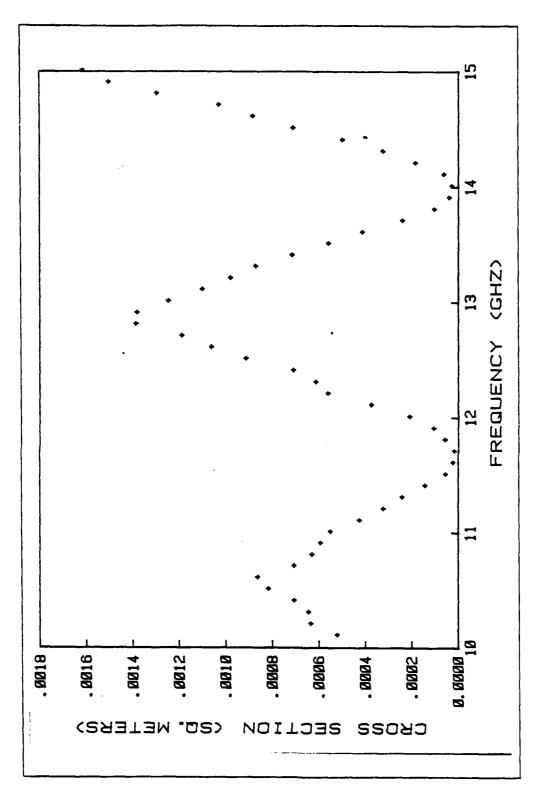


Figure 3.11 Cylinder 7: Cross-section

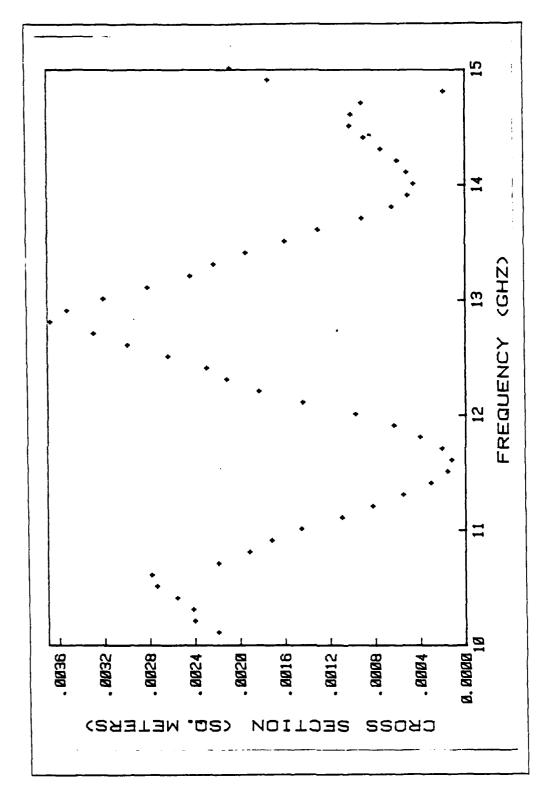


Figure 3.12 Cylinder 8 : Cross-section

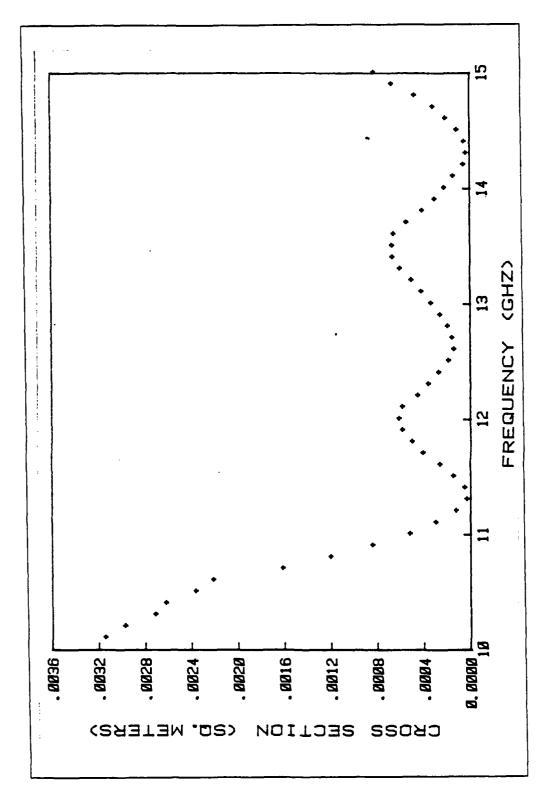


Figure 3.13 Cylinder 9: Cross-section

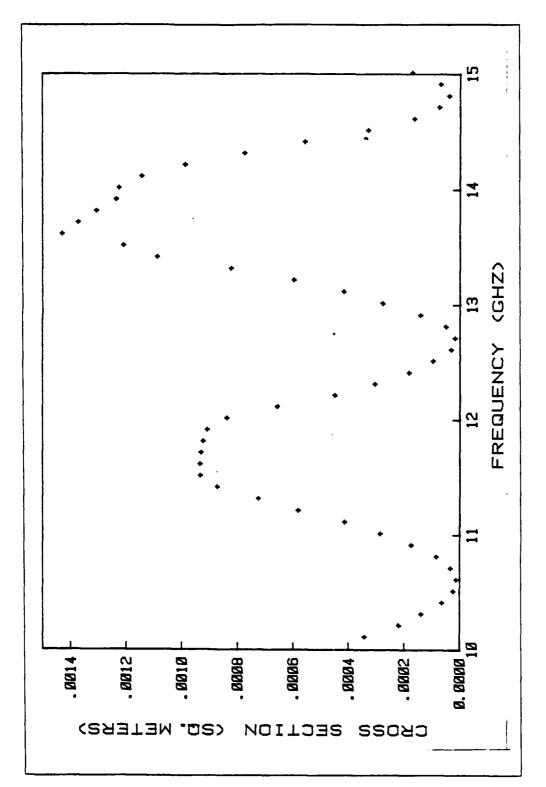


Figure 3.14 Cylinder 10 : Cross-section

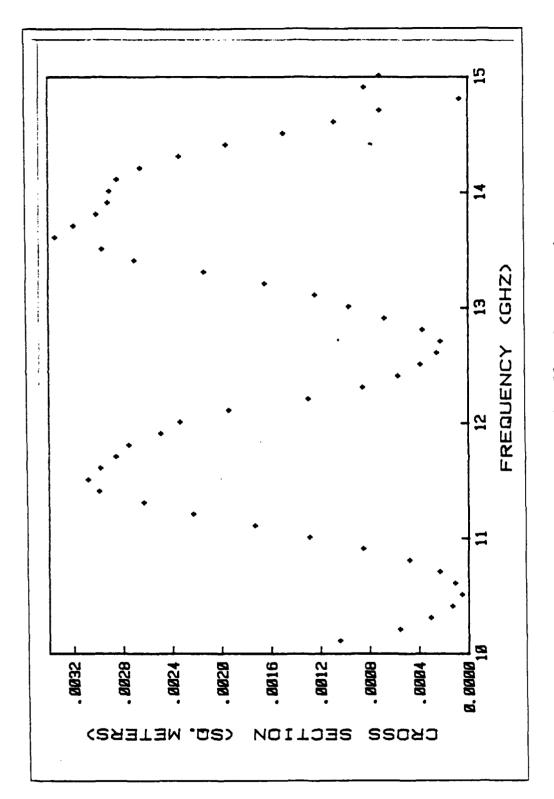


Figure 3.15 Cylinder 11: Cross-section

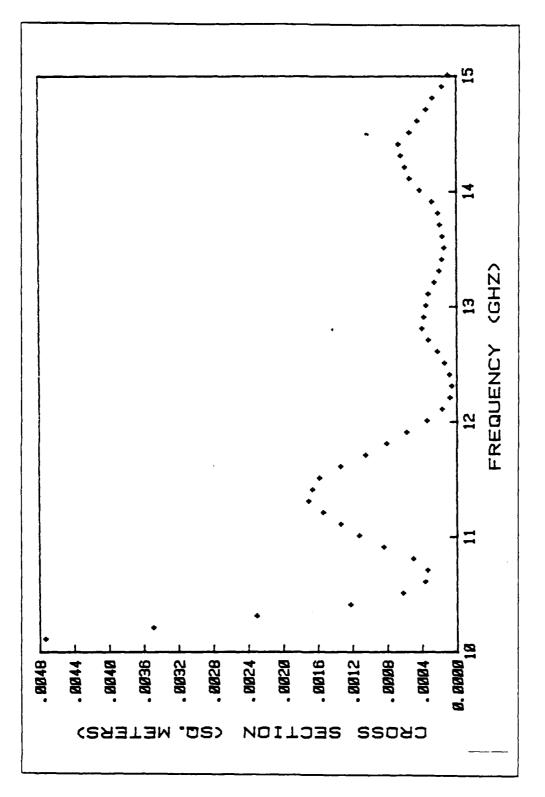


Figure 3.16 Cylinder 12: Cross-section

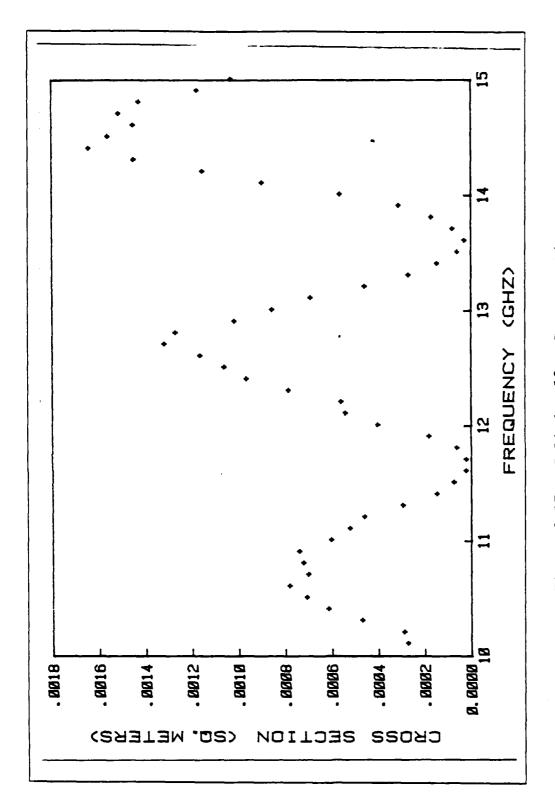


Figure 3.17 Cylinder 13 : Cross-section

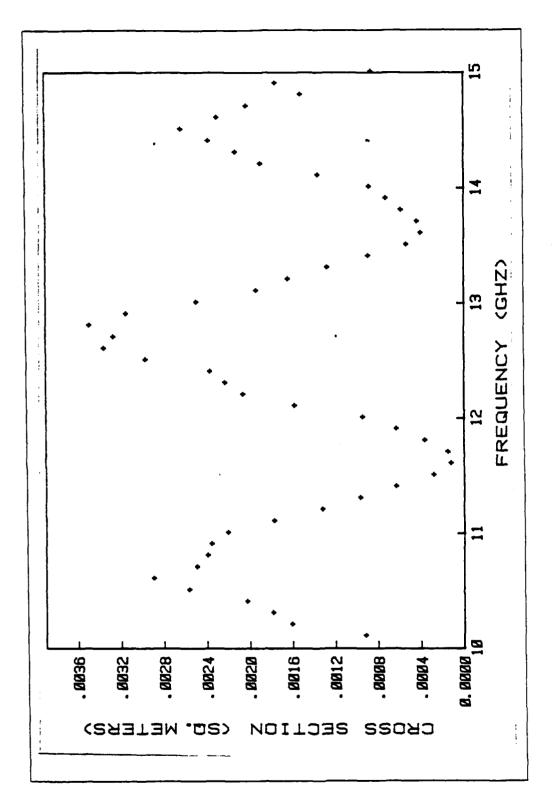


Figure 3.18 Cylinder 14: Cross-section

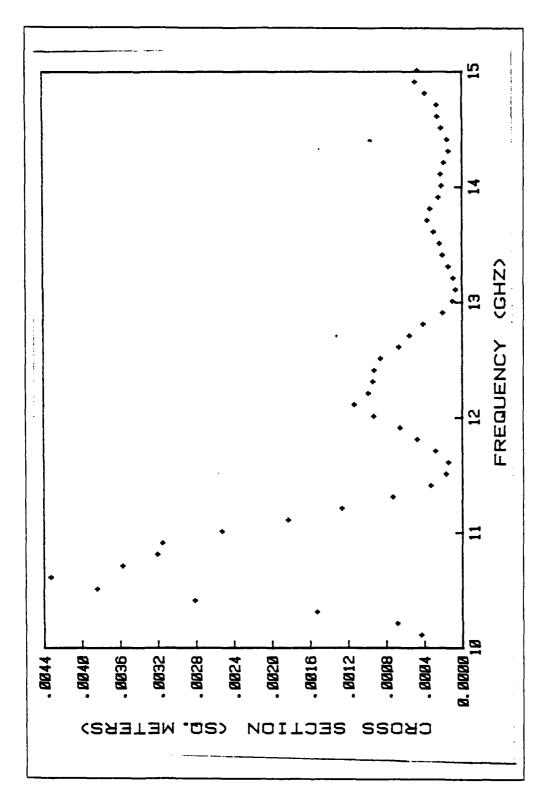


Figure 3.19 Cylinder 15: Cross-section

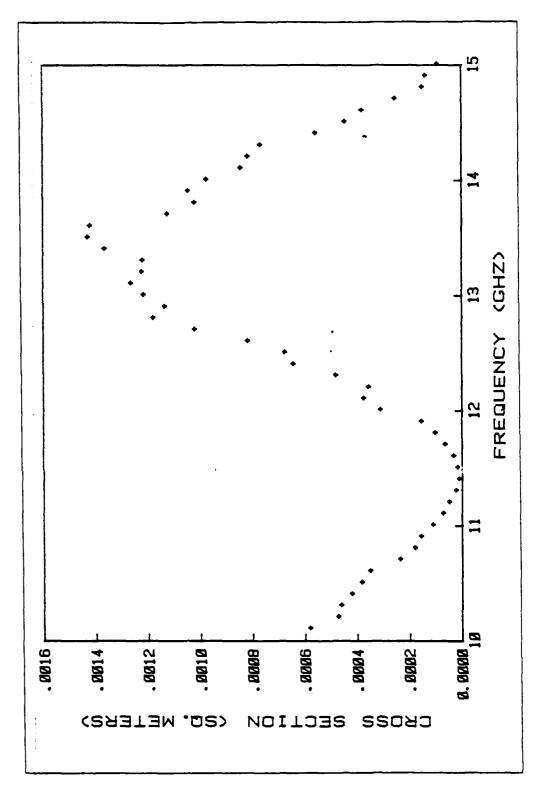


Figure 3.20 Cylinder 16: Cross-section

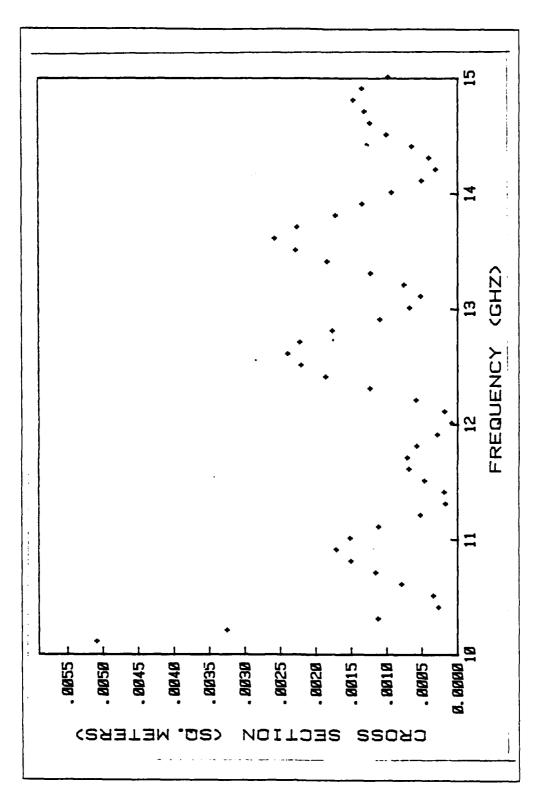


Figure 3.21 Cylinder 17: Cross-section

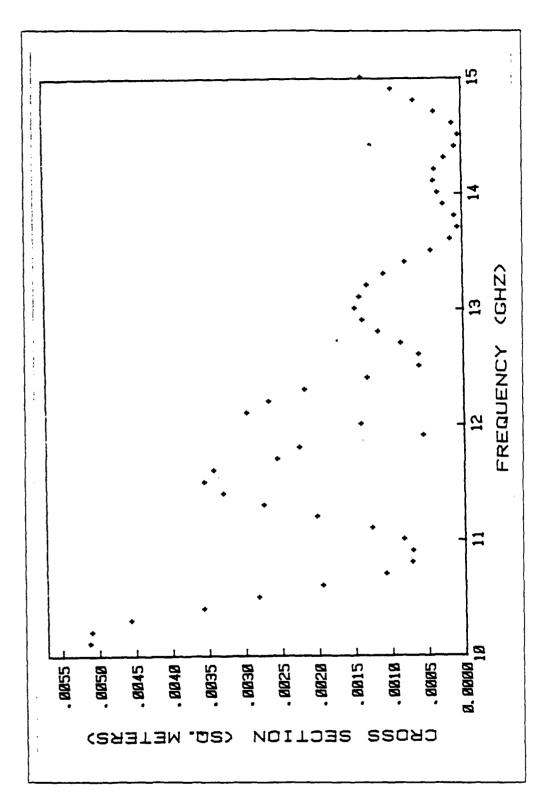


Figure 3.22 Cylinder 18 : Cross-section

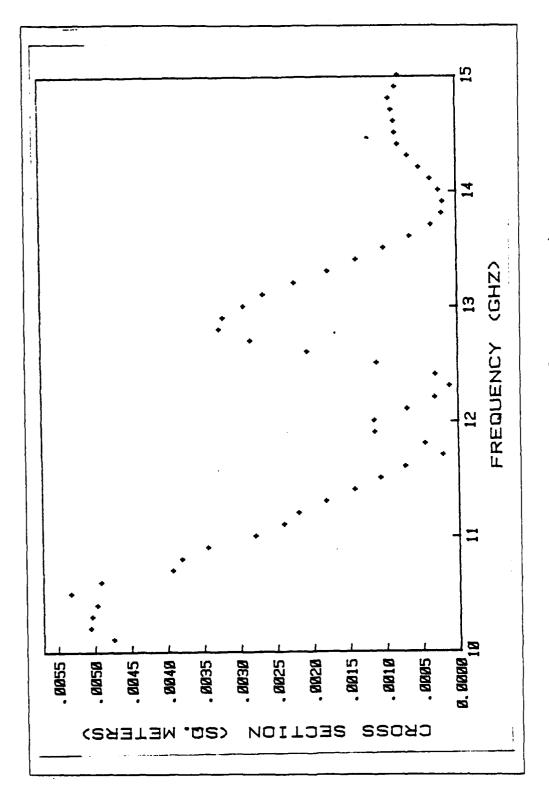


Figure 3.23 Cylinder 19: Cross-section

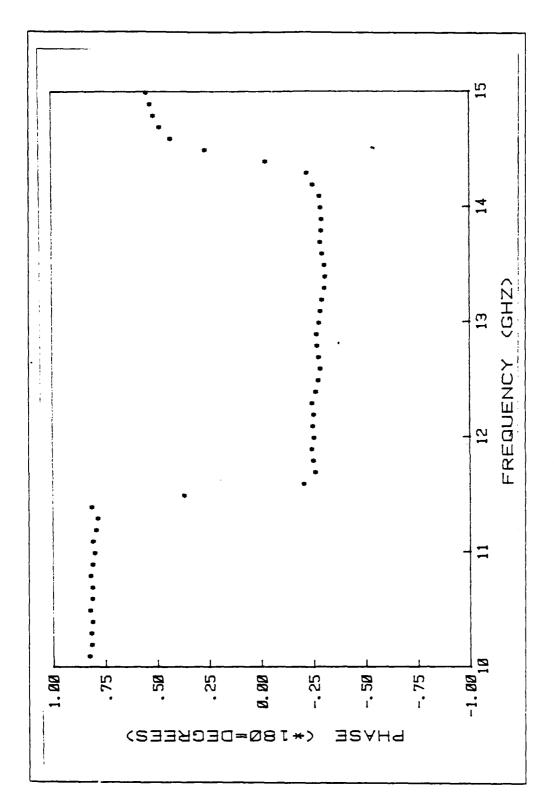


Figure 3.24 Cylinder 1: Phase

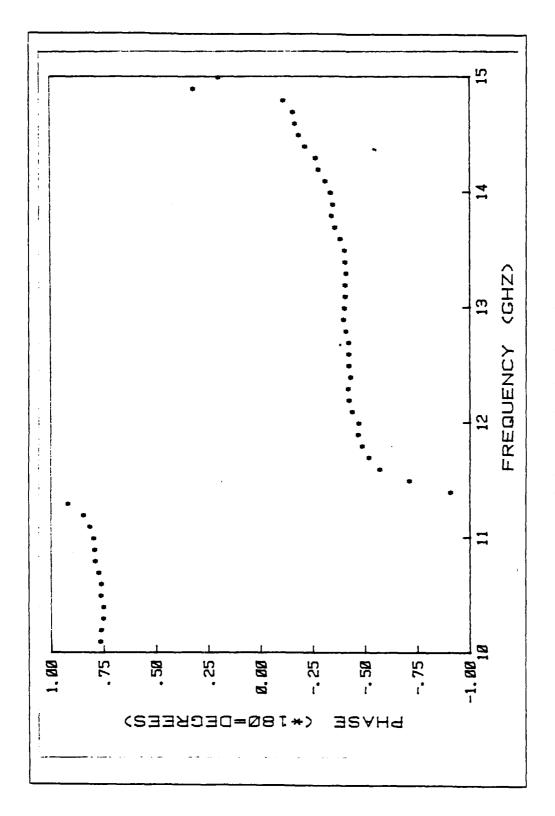


Figure 3.25 Cylinder 2: Phase

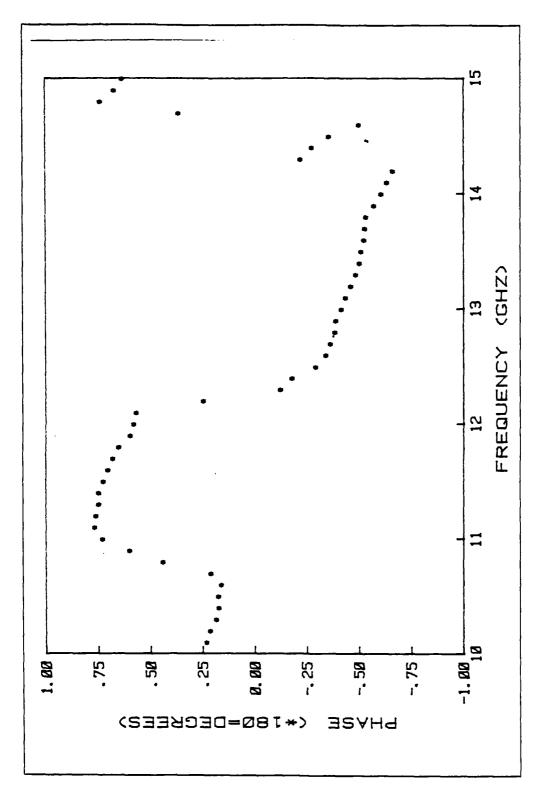


Figure 3.26 Cylinder 3: Phase

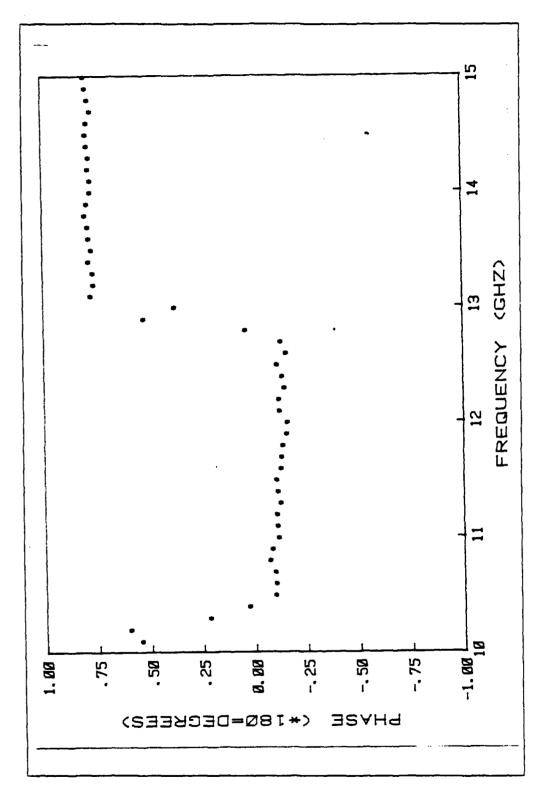


Figure 3.27 Cylinder 4: Phase

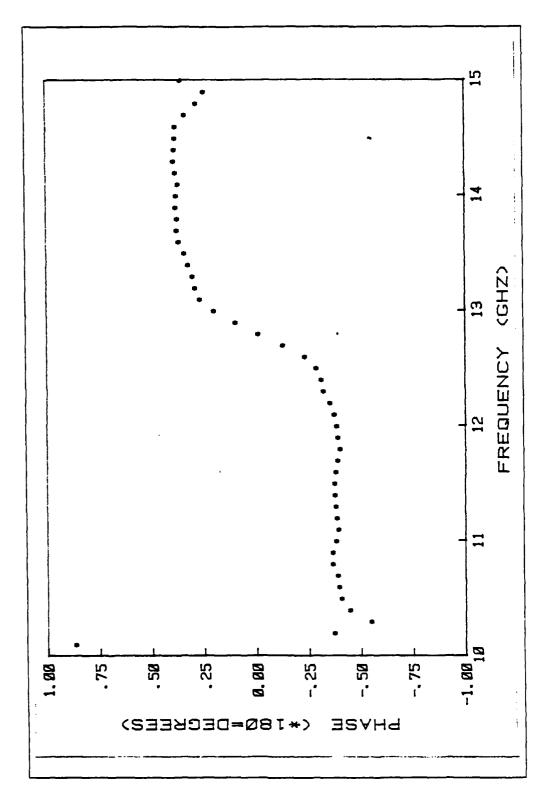


Figure 3.28 Cylinder 5: Phase

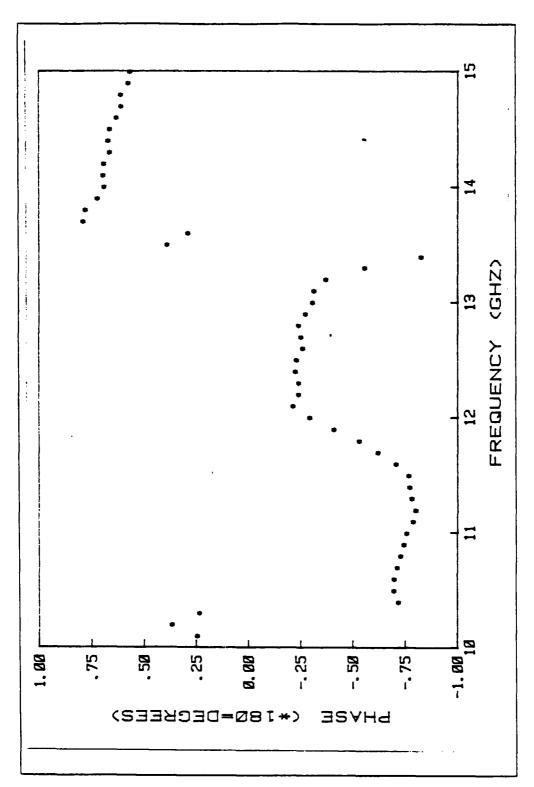


Figure 3.29

Cylinder 6: Phase

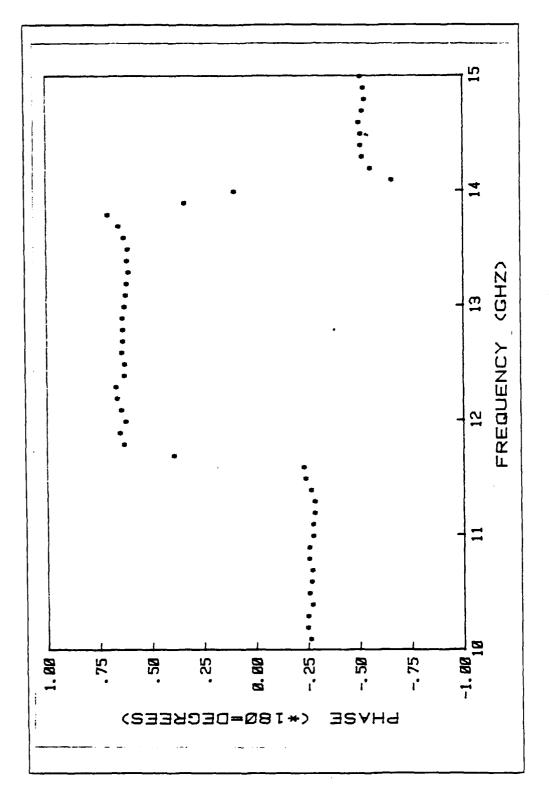


Figure 3.30 Cylinder 7: Phase

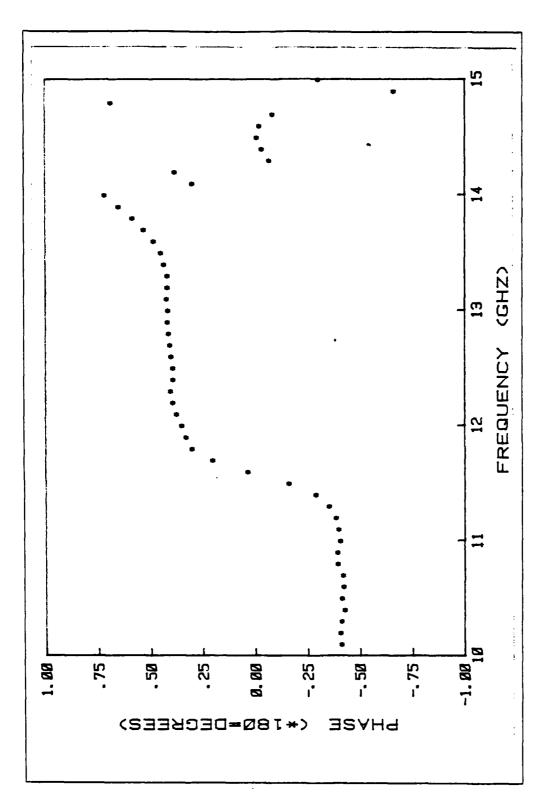


Figure 3.31 Cylinder 8: Phase

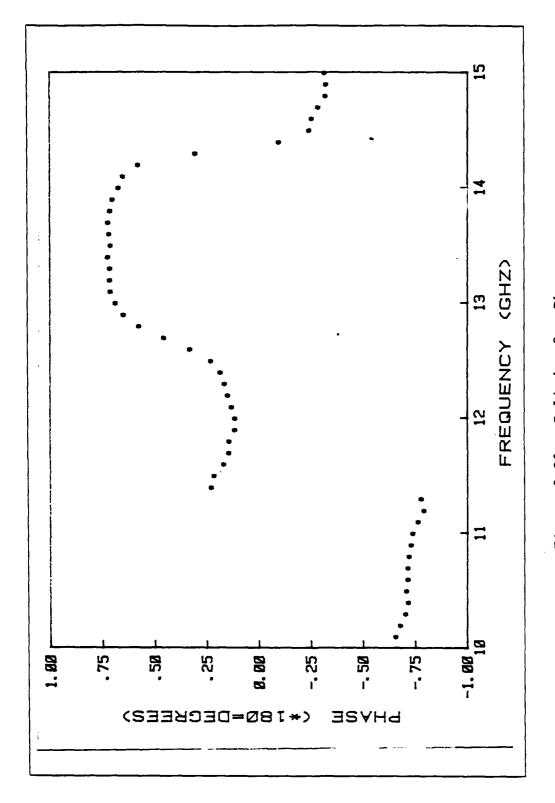


Figure 3.32 Cylinder 9: Phase

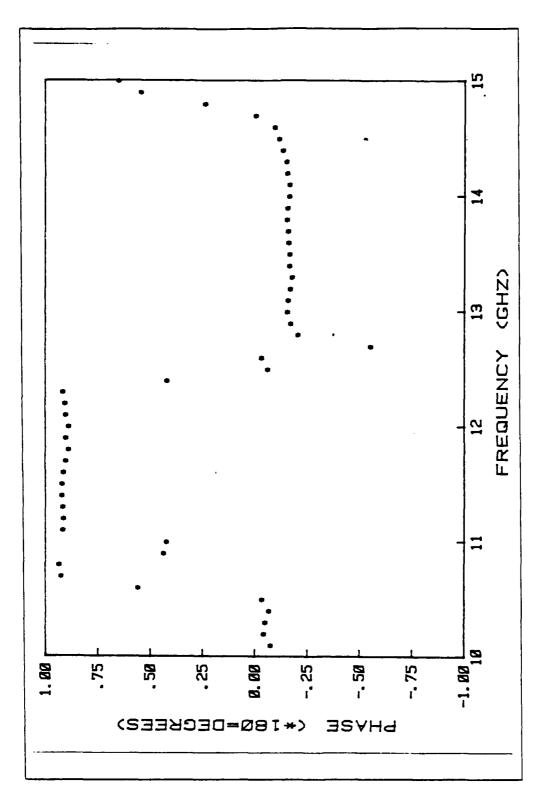


Figure 3.33 Cylinder 10: Phase

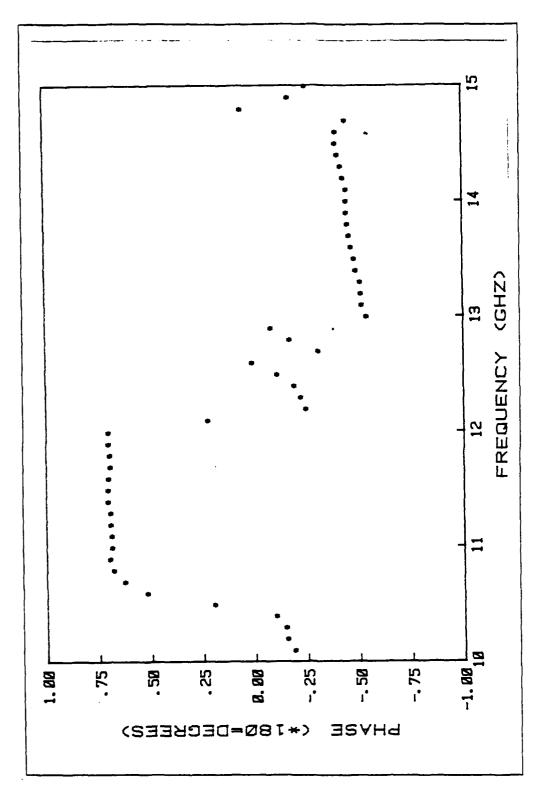


Figure 3.34 Cylinder 11: Phase

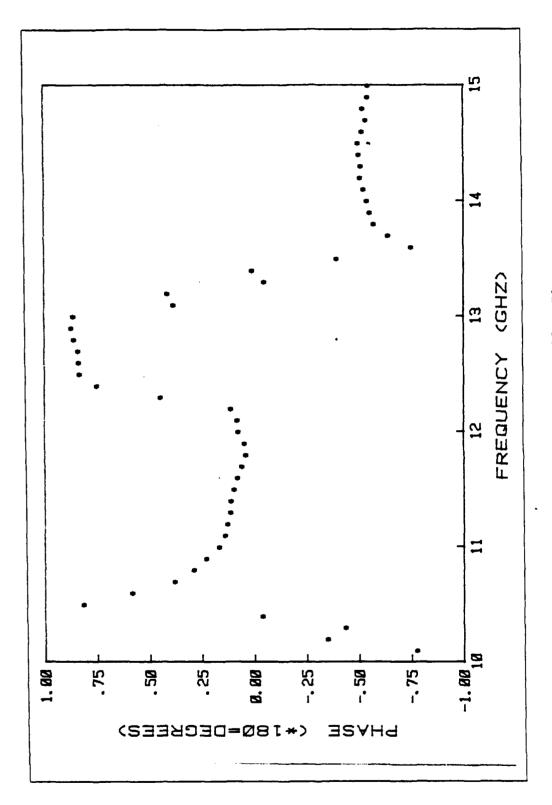


Figure 3.35 Cylinder 12: Phase

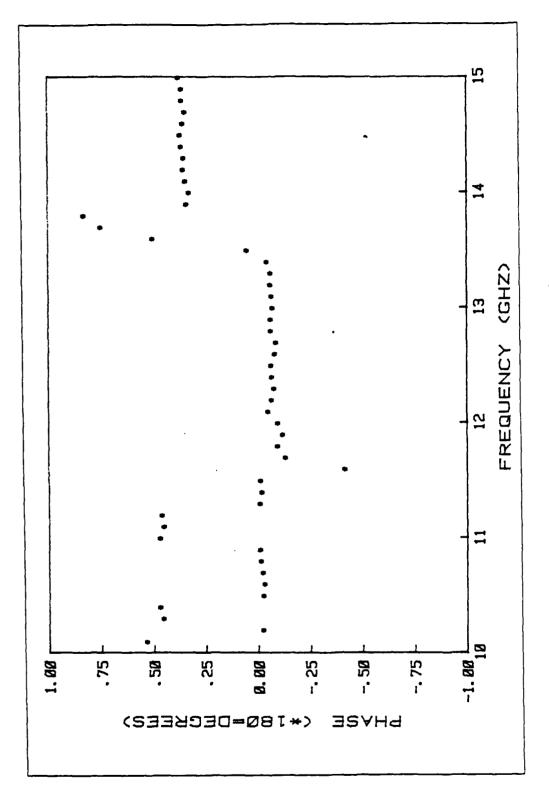


Figure 3.36 Cylinder 13: Phase

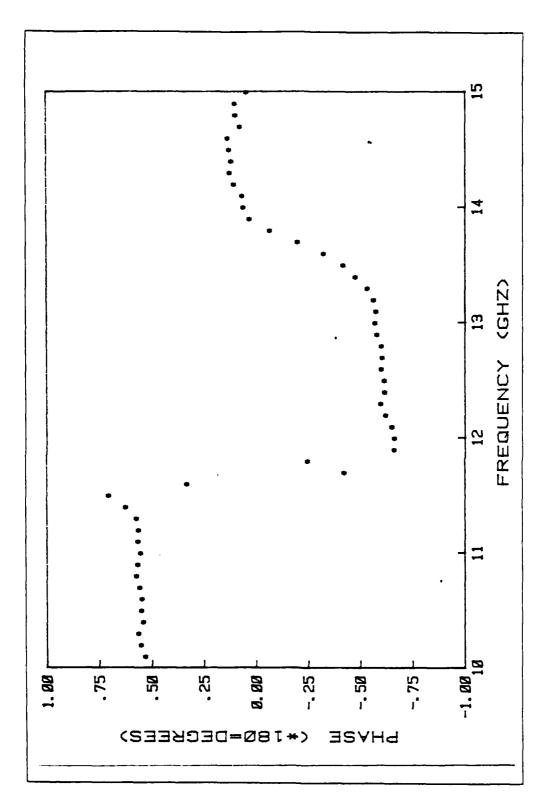
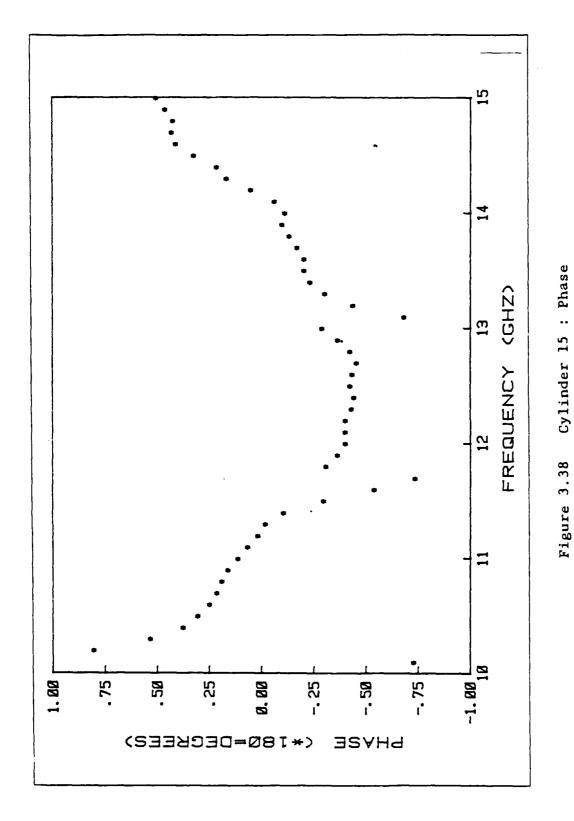


Figure 3.37 Cylinder 14: Phase



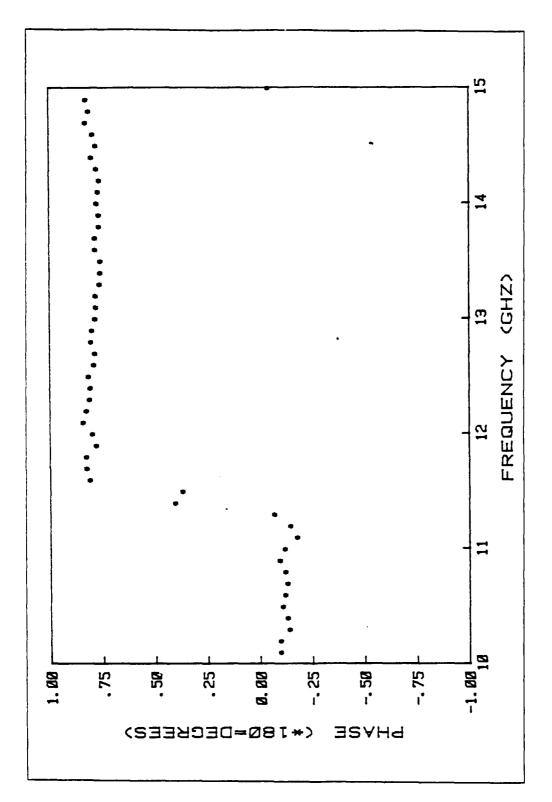


Figure 3.39 Cylinder 16: Phase

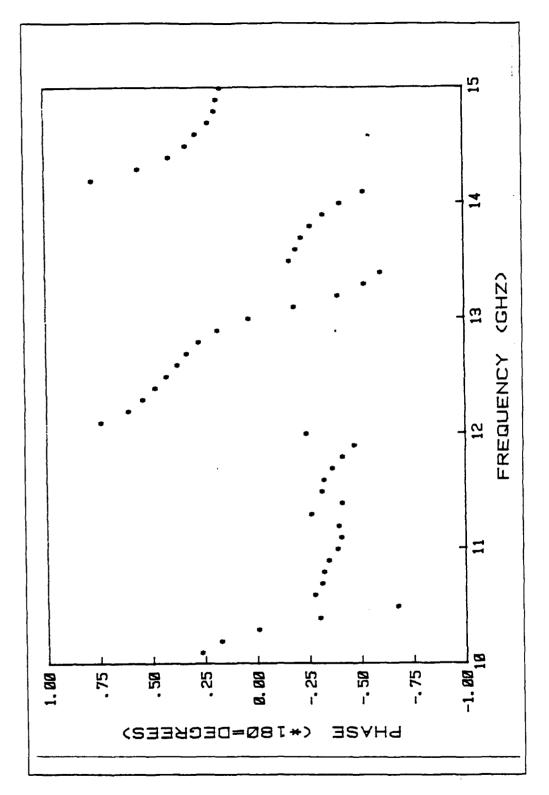


Figure 3.40 Cylinder 17: Phase

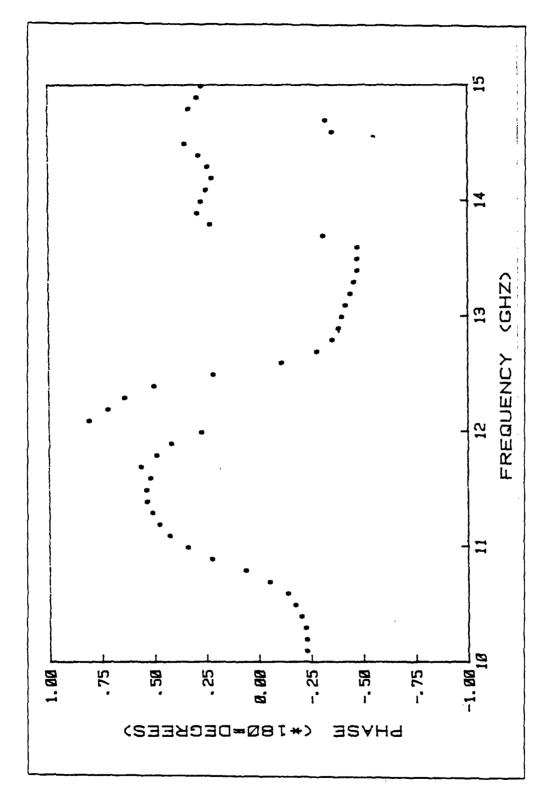


Figure 3.41 Cylinder 18: Phase

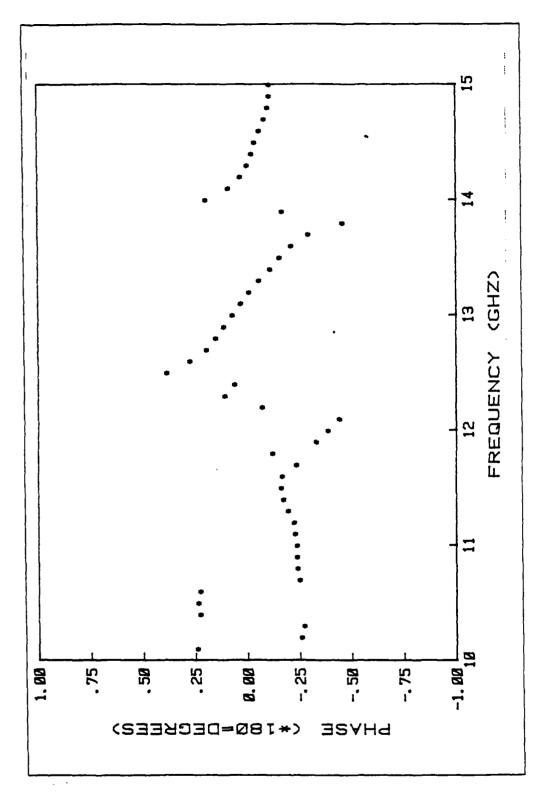


Figure 3.42 Cylinder 19: Phase

TABLE IV
CYLINDER 1 : MEASUREMENT DATA

FREQ	CRSEC	PHASE
(GHZ)	(SQR.METER)	(DEG /180)
10000000000000000000000000000000000000	79369834702159410265424669563777997757799745473253462 6777669854332100000000000000000000000000000000000	057815254911009211155452987966807699335745548403355117277761494949193022515245494040375777777777777455488027890498609609197410920729880949741092072988804039802080461756990472988804976335590472982222222222222233333333322200244912

TABLE V
CYLINDER 2 : MEASUREMENT DATA

FREQ (GHZ)	CRSEC (SQR METER)	PHASE (DEG./188)
10000000000000000000000000000000000000	3391110722759789504770854250449389830596343871090403221108722275978950477085425044938983059634438710904082223222222221111877533312222222211111877533312273333322222221111187753331227333332222222211110000000000000000	28203013382127833998770274626351931814017731403248 47568247501127833998770274626351931814017731403248 67543344594532188448548548279944005653837723373190 777777778830228308444444333333333333222211231 677777777777883997755544444444444444444444444444444444

TABLE VI
CYLINDER 3 : MEASUREMENT DATA

FREQ	CRSEC	PHASE
(GHZ)	(SQR.METER)	(DEG./180)
12745678988888888888888888888888888888888888	\$6046568086795728299974901862829150515465741867122 42873510117556776554329998992458011375665119875543111234 62011009090909090909099999111111115419875543111234 936090909090909099999999999999999999999	155175743785179553009800855050504918789479887506047766177755226668725953359492158571812421472655353594921585718124472655353595739555823948537474517537957895789579857492479386147357777777777777666555521113333444455555555555766223579661479667662235555555579644445555555555555555555555555

TABLE VII
CYLINDER 4 : MEASUREMENT DATA

FREQ (GHZ)	CRSEC (SQR.METER)	PHASE (DEG./180)
		0 1000617239259908599525837015744572524001231837272725 0 10006172239259908599533570157445725240012331837272725 0 1000617223925999204080533579990951703553540883072 0 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1
15.00	.00046	.78120

TABLE VIII
CYLINDER 5 : MEASUREMENT DATA

EBEA		OUGGE
FREQ (GHZ)	CRSEC (SWR.METER)	PHASE (DEG./180)
(dis)	7 W W W Y Y 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1	(040250
10.10	. ଉଷ୍ପଷ୍ଥ	. 85126
10.20	.00003 .00005 .00016	38443 5000
10.30 10.40	. 00000 00015	- 56296 - 46025
10.50	00033	41897
10.60	.00062	- 40914
10.70	. 00083	40267
10.30	.00116	- 37869
10.90 11.00	. 00151	37820 - 79541
11.10	.00197 00221	- 40625
11.20	00269	39541 40625 39817 39523
11.30	.00292	39523
11.40	.00302	- 39026 - 38968
11.50	.00307	38968 39398
11.60 11.70	.00497 • 00271	- 40377
11.80	99916 99916 999623 999623 999116 999116 999229 999397 999271 99274	- 41994
11.90	.00254	- 40500
12 00	.00234	- 40500 - 39843 - 38645
11.90 12.00 12.10 12.20	.00271 .00274 .00254 .00234 .00191 .00157	38645 36604
12.20	99137	- 33416
12.30 12.40	00113 .00078 .00059 .00045 .00030	36604 33416 32596 30133
12.50	. 00059	- 30133
12 60 12 70	.00045	- 30133 - 24710 - 14248
12.70	. 99935 . 999 <u>2</u> 5	
12.80 12.90 13.00	99931	08398
13.00	00031 .00043 .00057 .00074 .00107	. 18538
13.10 13.29	00057	25209
13.20	.00074	. 27611
13.30 13.40	. 00107 00146	. 40141 30947
13 50	.00146 .00184	32772
13 30 13 40 13 50 13 60	00210	.35150
13.70 13.80	.00224 .00240	36153
13.80 13.90	.00259	35898
14.00	. 90274	.300JJ 36405
14.10	00293	35618
14.20	00322	- 02427 023339 135209 12576117 22773472 336153 3364013 33664013 33664013 33664013 33664013
14.39	. 60312	.37422 .37218
14.40 14.50	.00300 .00276	.37218 36845
14.50	.00276	. 36789
14.70	00224	. 32210
14 30	.00225	26797
14.90	.00089	23036
15.09	. 00056	. 33964

TABLE IX
CYLINDER 6 : MEASUREMENT DATA

FREQ	CRSEC	PHASE
(GHZ)	(SOR.METER)	(DEG./180.
10000000000000000000000000000000000000	9598477647425685429574398529833183344132743486998889588998889521285743985298889889898989999999999999999999999	470825438636888844511440381212448988762988144532350171242426288844511440381212448988762988144532350171242426255189244327743813776998786289367689318342327777777788777776543822222333358377766666595555555555641948988624777666666555555555641948988627776666665555555556419489886288627776666665555555564194898862886277766666655555555641946766666555555556419467666666555555556419467666666555555556419467666666666666666666666666666666666

TABLE X
CYLINDER 7 : MEASUREMENT DATA

		-
FREQ	CRSEC	PHASE
(GHZ)		(DEG /180)
	(000):112:0:11	
		07544
10.10	.00051	27641
10.20	.00062	*26206
10.30	.00063	26263
10.40	99969	- 28534
	99997	- <u>4000</u> 4
10.50	. 00080	27050
10 60	.00085	- 28241
10.70	99969	- 28484 - 26957 - 27013
10.80	.00069 .00062	- 26057
10.00	. 99992	-,2030;
10.90	.00058	- 27013
11.00	.09054	28944
11.10	.00054 .00041	_ 20007
11.20	.00031 .00023 .00013	- 29767 - 29871
11.20	. 55531	20071
11.30	. 66653	7.470(1
11.30 11.40	. 00013	27982
11.50	. 00004 . 00001 . 00000	27982 25442 24695
11.60	99991	- 24695
11 00	00001	37470
11.60	. ଉତ୍ତରର	.3(4(8
11.70 11.80	. 00004	.37470 .61601
11.90	. 00009	. 53317
12 00	. 00009 . 00020	.60815
11.90 12.00 12.10 12.20 12.30	96976	.62896
12.10	.0003£ .00055 .00060	.04070 64055
12.29	. Genas	. 64955
12.30	. 00060	65435
12.40 12.50 12.50 12.20 12.00 13.20 13.30 13.50 13.70 13.70	.00070	.61442 .61292 .62563
12 50	00090	61292
12.00		60 6 67
12.50	. 99195	.04303
12.70	00118	.63120
12,80	.00137	.62072
19 90	.00137	. 62265
17 00	.00123	61241
13.00	. 99123	.61241 .60690
15.15	.00109	.00630
13.20	. ଡ୍ଡୁଡ଼୍କ	. 60 0 66
13.30	.00086 .00070	. 59325
17 40	99979	.59960
13.70 13.60	99955	59522
13.35	.00055	. 35322
13.60	.00040. .00023	.61393 .63828
13.70	.00023	. 63828
13.80 13.90 14.00	.00009	.69122 .32204
17 00	99997	72204
13.70	.00003 .00001	. 02207
14.00	. តីតិតិតិ T	.08397
14,19	.00005	67478
14.20	.00017	57053
14.30	00031	- 53141
	00001	- 52406
14.40		- J2465
14.50	. 00070	- 52495
14.60	.00087	- 51447
14.70	.00102	- 53291
14 80	00128	- 54296
14.90	.00149	53833
15.00	00160	52256

TABLE XI
CYLINDER 8 : MEASUREMENT DATA

FREQ	CRSEC	PHASE
(GHZ)	(SOR.METER)	(DEG./180)
10000000000000000000000000000000000000	7893169937987987851897166587918228339727989625 22257716993786444897166587918222222777169972788964444572798962 8888888888888888888888888888888888	1088051724970171035417190255925376510965824591864351053109727253118699727838393939317612091862459186434236269611862415626997288699788393939817612081233535359168347256444444444443333103333333333333444444444555672390006631444444444444444444444444444444444

TABLE XII
CYLINDER 9 : MEASUREMENT DATA

FREQ (GHZ)	CRSEC (SQR.METER)	PHASE (DEG./180)
100 100 100 100 100 100 100 100 100 100	259904999823701223886963356137731987443288992312989899999999999999999999999999	00427243767905112000112076792334427945051120001127723344631120005231005277777777777777777777777777777777777
15.00	. 00054 . 00050	- 32927

TABLE XIII
CYLINDER 10 : MEASUREMENT DATA

FREQ	CRSEC	PHASE
(GHZ)	(SQR.METER)	(DEG./180)
10000000000000000000000000000000000000	31351027670716222103354978214371918026921386525636609999999999999999999999999999999999	\$7248215978744412326515927845284372554334413108999835173594317164119977348226886752278785764164237548199754824885519878586472669647779886875227877857182858168284491936988888888446777878888888888888888888888

TABLE XIV
CYLINDER 11 : MEASUREMENT DATA

FREQ (GHZ)	CRSEC (SQR.METER)	PHASE (DEG./180)
10000000000000000000000000000000000000	.00102	3133911443279722633752021627465973273599910189276 20608971777972263359973273273599910189276 216131756681763179202169274659173273899109276 21613176681688779235997322845778899109276 2161317972266887792338457778999109276 21613179786888779238457788999109276 21613179786888779238478999109276 2161317978448597732735999109276 2161317978448597732735999109276 2161317978448597732735999109276 2161317978448597732735999109276 2161317978448597732735999109276 2161317978448597732735999109276 216131797847899109276 216131797847899109276 2161317978847899109276 2161317978847899109276 2161317978847899109276 2161317978847899109276 2161317978847899109276 2161317978847899109276 2161317978847899109276 2161317978847899109278999109199276 2161317978847899109278999109199276 21613179788478991099109199276 2161317978847899109199276 2161317978847899109199276
15.00	. 99969	26882

TABLE XV
CYLINDER 12 : MEASUREMENT DATA

FREQ	CRSEC	PHASE
(GHZ)	(SQR.METER)	(DEG./180)
11.60 11.80 11.80 11.80 11.90 12.20 12.20 12.30 12.80	96799306090072401741352519964292631363530503011436 96799306090072401741352519964292631363535503011436 9609000000000000000000000000000000000	6693002501067820057732801703993399582695972715922206815342501067820657732801722921135511397572715922206815324250106867550222947607229211355113946881779298642364236676933584567476886747635667476356747635555555555555555555

TABLE XVI
CYLINDER 13 : MEASUREMENT DATA

FREQ	CRSEC	PHASE
(GHZ)	(SQR.METER)	(BEG./180)
10 10	.00026	.52011
10 10 10.20	.00026	03685
10.20	.00025	44106
10 40	иииби	45625
10.50	.00060 .00070	- 03695
10.60	.00077	- 04266
10.70	.00069	03491
10.80	.00071	- 02670
10.90	.00073	- 02308
11.00	.00059	45700
11.10	.00051	. 43799
11.20	.00045	.44603 02343
11.30	.00028	- 03048 - 03048
11.40 11.50	00002	02486
11.50	. 90900 00001	42811
11.70	00001	- 14594
11.80	คัดคัด5	- 10856
11.90	00017	- 14594 - 19856 - 13184
12.00	.00039	- 10895
12.10	.00053	10895 06273
12.10 12.20	. 00055	- 07983
12.30	00070 00070 00077 000077 000051 000051 000001 000001 000007 00000 00000 00000 00000 00000 00000 0000	07983 09254 08040 07835 09568 10050
12.40 12.50	. 00095	08040
12.50	. 00105	07833 00560
12.60 12.70 12.80 12.90 13.00		-,07068 - 10050
12.70		- 10000 - 07724
12.00	99199	- 07662
13 00	00084	- 07784 - 07662 - 08504
13.10	.00126 .00100 .00084 .00068 .00044	- 07916 - 07499 - 07608
13.20 13.30 13.40	.00044 .00026	07499
13.30	.00026	07608
13 40	.00013	- 45264
13.50	.00005	.03834
13.50 13.50 13.60 13.70 13.80	99995 99992 99997 99916 99939	.48835 77770
13.70	. ଅପ୍ରଥମ : ଜନନୀ କ	21574
13.00	00010	32548
14.00	00055	03834 43856 73730 81574 32548 31314 33014
14.10	00089	33014
14.20		
14.30	.00144	. 33673
14.40	. 99163	.34956
14.50	00155	.35459
14.60	.00144	.34276
14.79	.00150	33157
14.80	.00141	. 34664 . 3468 3
14.90	00116	36290
15.00	.00102	. ಎರ೭೨೮

TABLE XVII
CYLINDER 14 : MEASUREMENT DATA

FREQ (GHZ)	CRSEC (SQR.METER)	PHASE (DEG./180)
10.10 10.20 10.30 10.40 10.50 10.60 10.70 10.80 11.00 11.10 11.20	9860477738594169240153044344141618957000000000000000000000000000000000000	7 754977039683322088024725443322031040680516575297914 7788995009745838024725443322031040680516575229797 7 55423344899555344896776523048805727563111193780222737 8 5555544889677765321104899771239330141098330095 7 7 7 7 7 7 7 7 7 7 7 7 7 7 7 7 7 7 7
15.00	. 00083	.02986

TABLE XVIII
CYLINDER 15 : MEASUREMENT DATA

_		
FREQ (GHZ)	CRSEC (SQR.METER)	PHASE (DEG./180)
Z G 123456789000000000000000000000000000000000000		8430031450467605280305094218698051264988516476325 29926603147632763190242405994218698051264985334625444733 77815833974631902424059924950640927995334625444738 21114900311049277156389627995334427733244444444443364514773344444444443364514773344444444443364322111127354444444444336451477334444444444336451477334444444443364514474444444443364432444444444444444444444
14.89 14.90 15.00	. 99935 . 99946 . 99943	.44520 .48910
	· · -	

TABLE XIX
CYLINDER 16 : MEASUREMENT DATA

TABLE XX

CYLINDER 17 : MEASUREMENT DATA

FREQ (GHZ)	CRSEC (SQR.METER)	PHASE (DEG./180)
10000000000000000000000000000000000000	418284116777834424734444911658253818944288996768697382828284167778344247344449116582538189442888967686977383889888888888888881822218384428889676869773838888888888888888888888888	113884440611136769599842906555226372962209969932955178 16789449061113667699599842906555226372962299555178 2152149496797489952004466550533853809981677889735444 21521359142137288423772844255555555555555644235334426644 21521494235334423333442755444333216240318135353998221664 114526442533344275544433321622234234357533221644 114526442333442333442755444333216246445142233435333221644

TABLE XXI
CYLINDER 18 : MEASUREMENT DATA

TABLE XXII
CYLINDER 19 : MEASUREMENT DATA

FREQ (GHZ)	CRSEC (SQR.METER)	PHASE (DEG./180)
123456789000000000000000000000000000000000000	963381387927051811262237671234774592271953051831 00643719066729753352253961235581334774592271953051831 006437190667297353352261255811347731000223322000000000000000000000000000	7 38666567201861204881788440797521000212851899910767467 8 42869268830530015300212854455964831569832593532248571 7 2222110002344522047975028943562889733837667419 8 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1
14.90 15.00	.00091 .00132	25345

IV. SUMMARY AND CONCLUSIONS

A. ANALYSIS OF MEASURED DATA

The back scattering cross section of a tubular cylinder depends upon several parameters. To simplify the problem, many of them were kept constant: The polarization of the receiving and transmitting antennas were kept constant during all measurements, all the cylinders were from the same material and the aspect angle and tilt of the different cylinders relative to the antennas plane were unchanged.

The parameters that were changed and their effects on the cross section were the subjects of this study were:

- 1. The cylinder length. (2h)
- 2. The cylinder diameter. (2a)
- 3. The transmitting frequency.

The cross section of a tubular cylinder can be written as a function of three parameters as shown in equsion 4.1

$$\sigma = F_1(a,h,f) = F_2(a,h,k)$$
 (eqn 4.1)

where:

$$k = 2\pi/\lambda = 2\pi f/c$$

$$c = 3 \times 10^8 \text{m/sec}$$

From the theory described in Chapter the following relations can be obtained:

(eqn 4.2)

$$\frac{\sigma}{\pi a^2} = F_3(l_1, l_2) = F_4(l_1, \frac{l_2}{l_1}) = F_s(ka, \frac{h}{a})$$

where:

$$l_1 = ka$$

$$l_2 = kh$$

The measurements taken can be devided into two steps. In the first step, 19 cylinders with different lengths and diameters were placed in the anechoic chamber, and plots of cross section and phase versus frequency were obtained on each of the cylinders. (figures 3.5 through 3.42) Except for gathering data and viewing the resonance frequencies, there is nothing much to say on the results because too many parameters were changed and the resulting analysis became too difficult and complicated.

The second step of measurements was based on the conclusion from the theory (equation 4.2).

Only eight scaled cylinders were put into the anechoic chamber, four of the cylinders had a constant ratio of length to diameter equals 4, (h/a=4), and four cylinders had a constant ratio of length to diameter equals 6, (h/a=6).

An overall plot with expanded frequency range was achieved by using those scaled cylinders.

The overall plots can be compared to the theoretical plots obtained by Prof.H.M.Lee at the Naval Postgraduate School.

Both the theoretical and experimental plots of cross section/area and phase versus ka are shown in Figures 4.1 through 4.8 and the overlapping ranges of the scaled cylinders in ka can be noticed from Tables XXIII through XXX The theoretical and experimental curves can be seen on the same graph for the two sets of scaled cylinders in Figures 4.9 and 4.10

Comparison of the experimental results to the theoretical plots show agreement away from the cutoff frequencies of the H_{11} circular waveguide mode at ka=1.8415.

In the theoretical calculations an infinitesimal wall thickness is assumed for the tubular cylinder, but practically the cylinders used in the measurements had some thickness (Table III). The outer diameter of the cylinder is used for 2a in the computations while the H_{11} mode cutoff frequency depends on the inner diameter of the cylinder. This fact caused the deviation in the cross section plots between the theoretical curve and the experimental curve near the cutoff frequency.

Problems occur in the phase plot because the averaging procedure did not properly take care of the phase shifts with values near \pm 180 degrees. Because of the complicated behavior of the phase shift near the H $_{11}$ cutoff at ka=1.8415 and near ka=2.4046, no conclusion can be made related to the actual behavior of the phase shift curve.

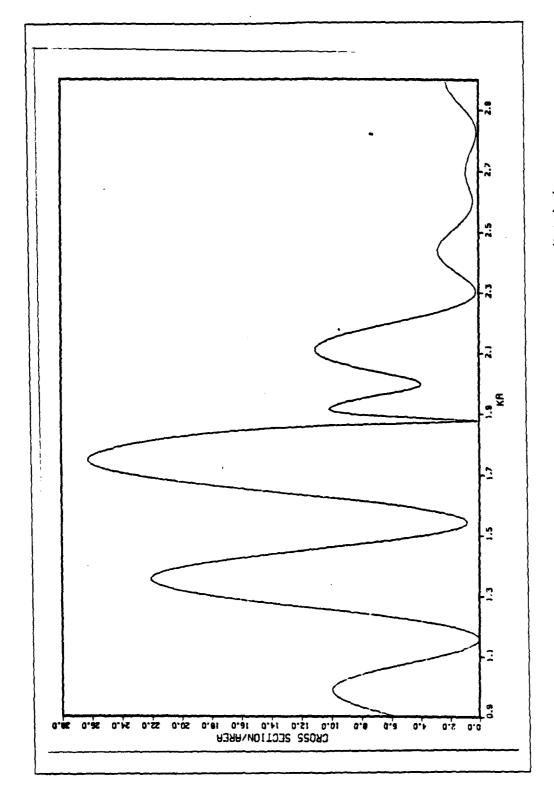


Figure 4.1 Theory Cross section Curve, (h/a)=4

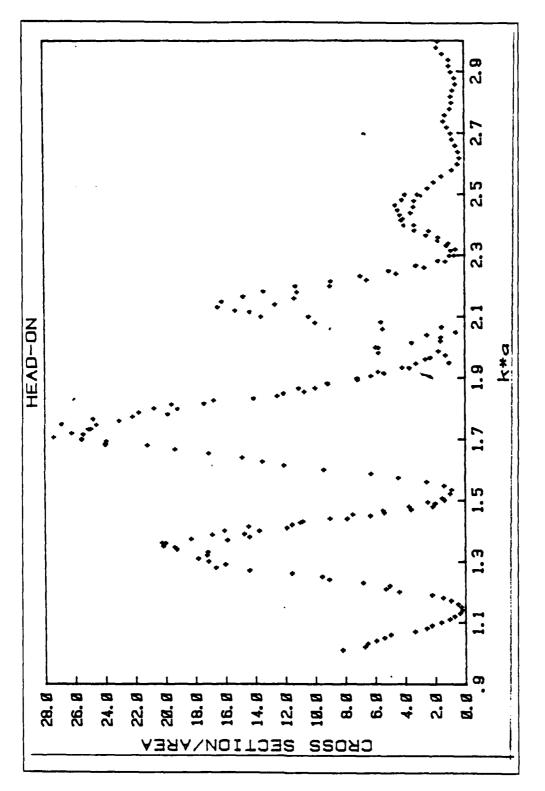


Figure 4.2 Experimental Cross section Curve, (h/a)=4

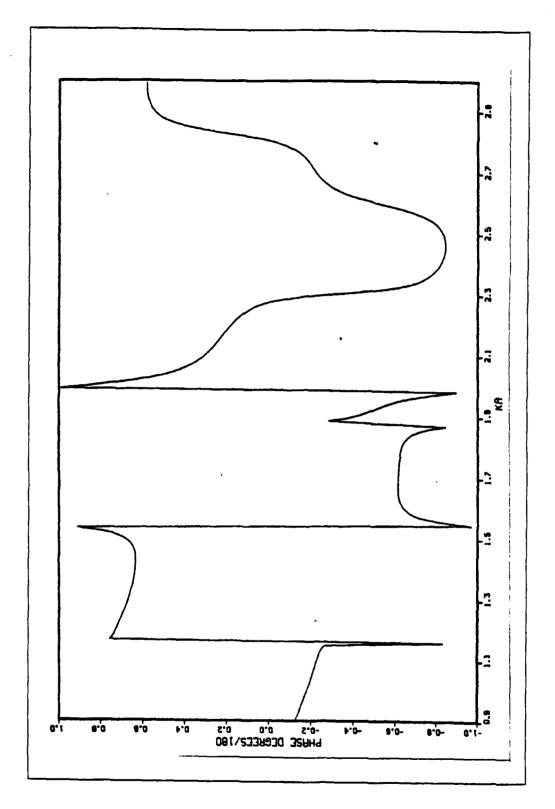
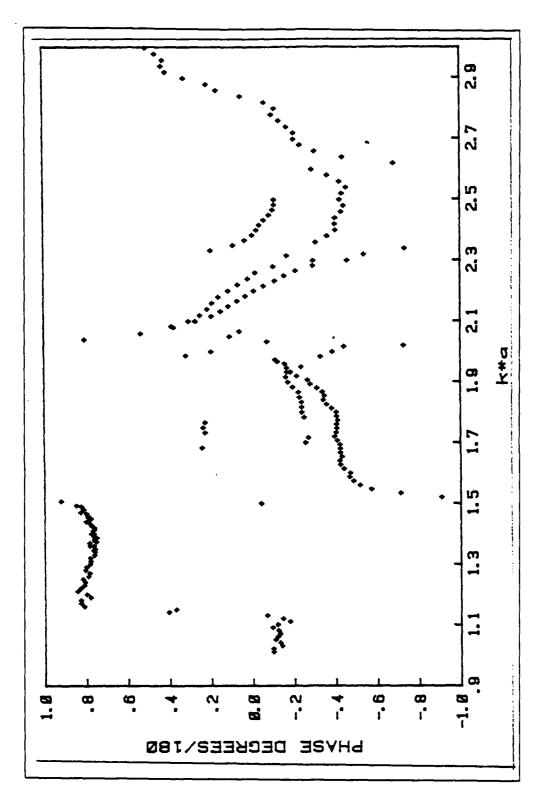


Figure 4.3 Theory Phase Curve, (h/a)=4



igure 4.4 Experimental Phase Curve, (h/a)=4

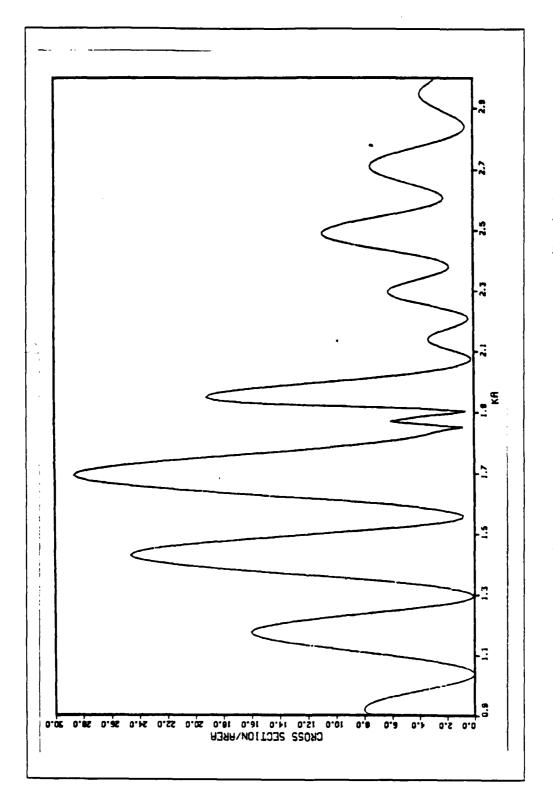


Figure 4.5 Theory Cross section Curve, (h/a)=6

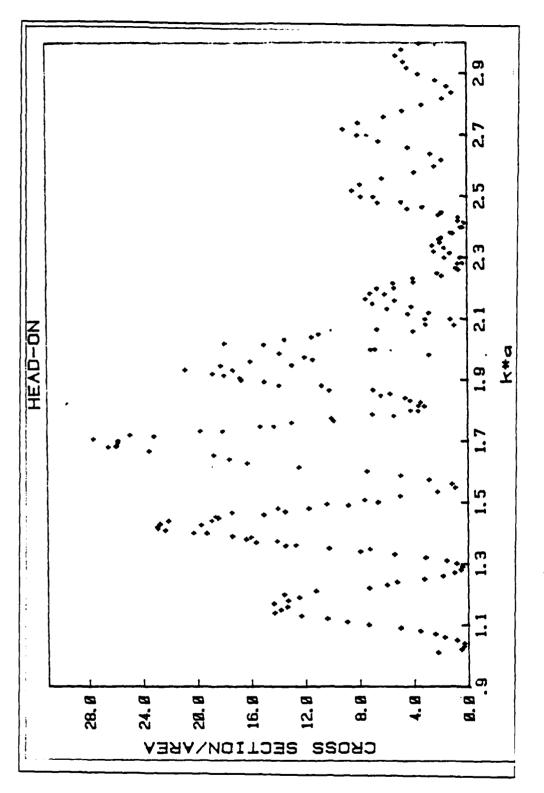
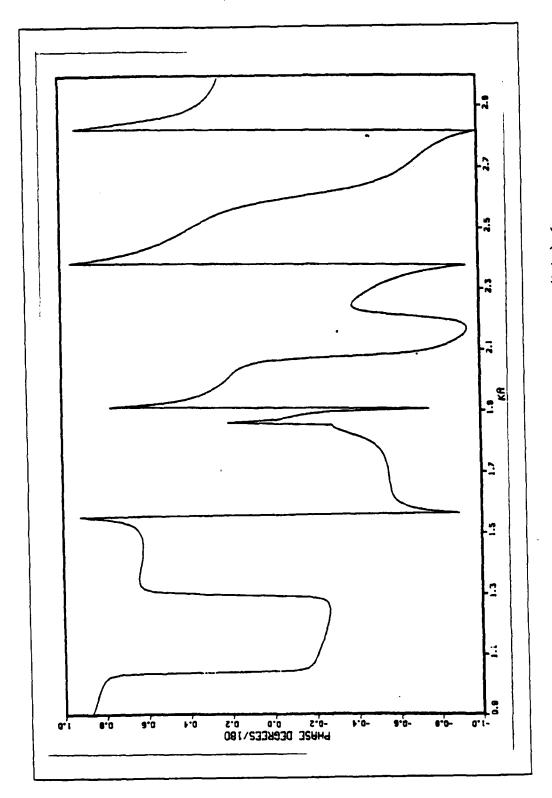


Figure 4.6 Experimental Cross section Curve, (h/a)=6



igure 4.7 Theory Phase Curve, (h/a)=6

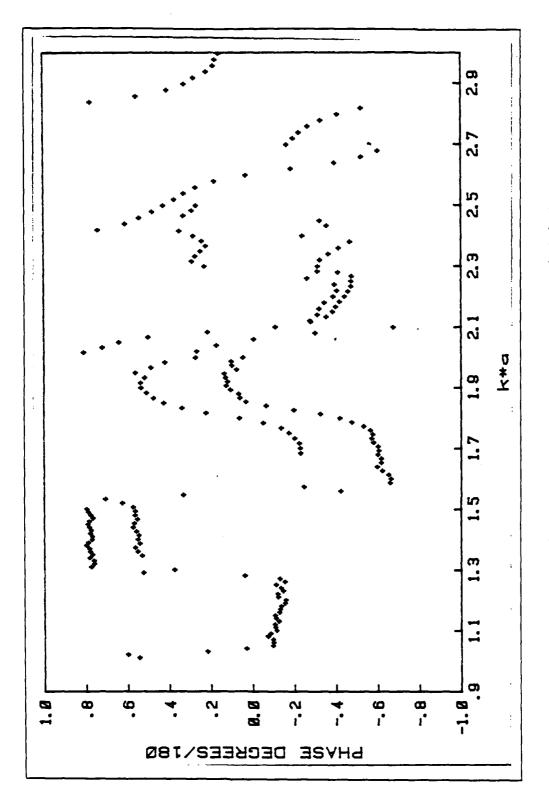


Figure 4.8 Experimental Phase Curve, (h/a)=6

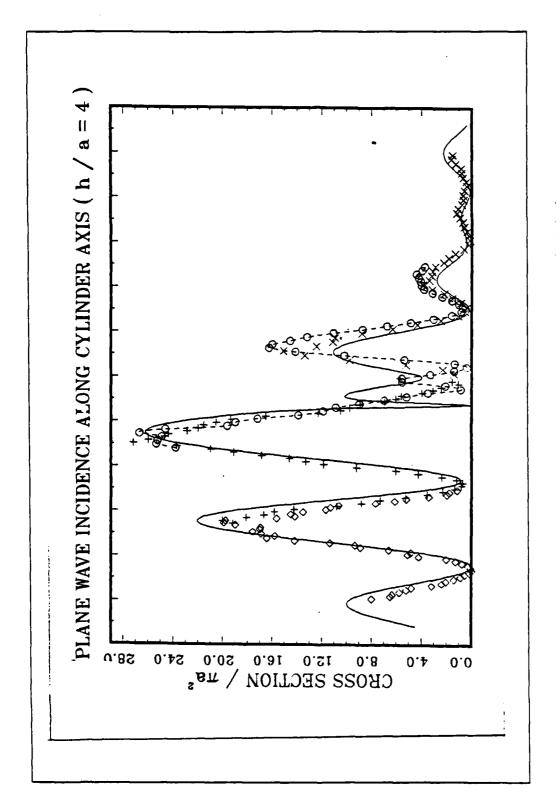
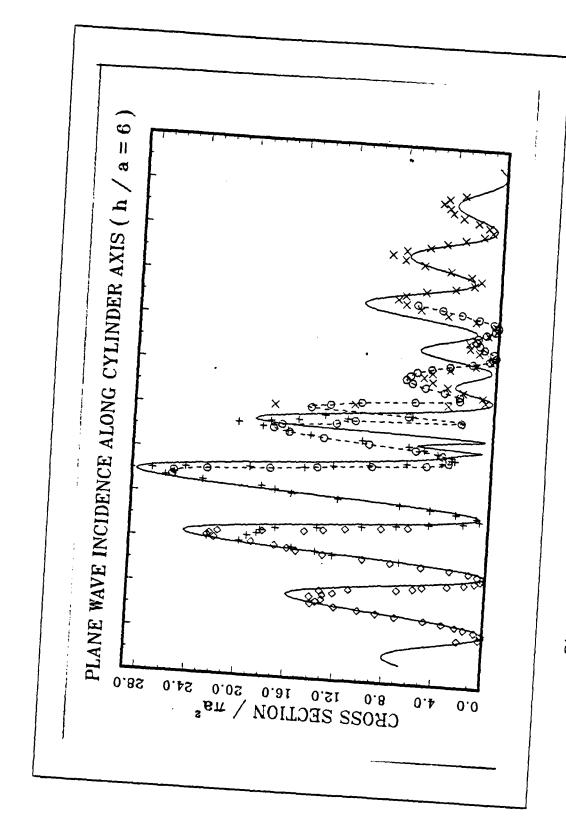


Figure 4.9 Summary Cross section Curve, (h/a)=4



igure 4.10 Summary Cross section Curve (h/a)=4

TABLE XXIII
CYL16.4:1 RATIO

(K * A)	CRS/AREA	PHASE
1234567898123456789812345678981234567898123456789812345678981234456789	H ## ## ## ## ## ## ## ## ## ## ## ## ##	39213207169359027311229803876896846766314848036019 212338447169359027311122980388768946766314848036019 21136235430261977759111195439632014094897710697719697 2111111111111111111111111111111111111
1.50	1.10104	A CONTRACT CONTRACT

TABLE XXIV
CYL.2,4:1 RATIO

(K #A)	CRSZAREA	PHASE
46780124568902346780124568902346780124568902346780124568902346780	906012888197744875633379115842852782184961615210342874558831474633379115842852188496161521033428085331862324423633308134428527822088623748153208282828188775332208961615498535844669113469905753844221109866533221152469113469903775442211098665332115	28203030133821278339987702746263519318140177314032486790364750011517321885548626351981801777314032486790469463631517321885548627904400565383772337310077777777888999755544444444444444433333333222211123344444444

TABLE XXV
CYL.18,4:1 RATIO

				211225
	(K #A)	CRS/P	REA	PHASE
1.	677777788888888999999999999999999999999	22222221111111111111111111111111111111	7892521701353602235081060153212047763332502778	756477783777498355897064875539022638211100994583583797727812812123552683770648755355764181999945835781222222222222222222222222222222222222

TABLE XXVI
CYL.15,4:1 RATIO

(K*A)	CRS/AREA	PHASE
1357913579135791357913579135791357913579	55844992455274250518423837481828393462445392603477183812736427383744931828374803462445392603477583356427364428064287288312756685741696990387273041211086421	8430031450467605280305094218698051264885164763250071450467605280305094218698051264885164763250077316094874504649327775669924245700564093277333444457889703640322211127359558211149509311554277114457889703643222111127351995582977756432221111273544848444444433643222111127334484848484848484848484848484848484848

TABLE XXVII
CYL.4,6:1 RATIO

(K *A)	CRSZAREA	PHASE
12345678901234567890123456789012345678901234567890 111111111111111111111111111111111111	2 12347899586686997779939591398247761885392882628262939586869977739866859913982477618853928826253773388689913982477614892882128137769248289313913986287124776128662577693121311775531 22221843131175531 22222184311186	100061723925908599525837015744572524001231837272520 5820617239926995258370157445725231837272520 11147399231460953357015744674609123122445692531645399525534683075520 11224456925316095170377777777777777777777777777777777777

TABLE XXVIII
CYL.14,6:1 RATIO

(K ≭A)	CRS/AREA	PHASE
4678012456890234678012456890234678012456890234678012456890234678012456890234678012456890234678012456890234678012456890234678012456890234678012456890234678012456890234678012456890234678012456890	102338104883023935268437597408819483775113502481732231188742233435268735934352687359343526873237456450687323752687351597408819483795564506873237555597564763486372221111111111111111111111111111111111	754977039683322080247254332208310406805165752979146 7782997468332208024725433220310406805152979146 77828555344509077801577633130496805161155752979146 7782331331314557531111286357 778231313146599775531111286357 778231313146591131268357 77823131465775531111286892

TABLE XXIX
CYL.19,6:1 RATIO

	_	
 (K #A)	CRSZAREA	PHASE
11111111111111111111111111111111111111	3908 448 592 447 447 441 441 441 441 441 441 441 441 441 441 441 441 442 443 445 445 445 445 445 445 445 445 445 445 445 445 445 445 445	38665567201861204881784407975210021285189910767467542818185648830940817844407975210021285189910767467522218564883094081784440797521002128518991076746754459648833628897335322485719444484576889733353224857194448336288973335322485753175

TABLE XXX
CYL.17,6:1 RATIO

(K*A)	CRSZAREA	PHASE
1357913579135791357913579135791357913579	1713 123555531 1221 146787632124678755431 12344543	111388844406111367695998429065522637296229955178 2151888444061113676959984290655226372962299735973544 216233597748536656553388099209969973544 21724237973804425653338096297359973544 21728842239973544 2172842238973599224 2184238917359998229987 2283423897359998229987 24556181222334256387 24561812223342363788

B. CONCLUSIONS AND RECOMENDATIONS

The results showed agreement between the theory and the experimental data obtained. Some deviations between the theoretical and experimental plots could be explained.

The problem caused by averaging as observed in the phase plot can be overcomed by refining the program in the following way: If the two values to be averaged deviate more than a pre-determined value, say, 200 degrees, both will be converted to a positive value before the average is taken. The averaged value is then adjusted so that it lies between -180 and +180 degrees. A new set of measurements with targets having length to inner diameter ratios of 4 and 6 is planned. This refinment of averaging procedure will be adopted.

For the next step in the comparative study of target back scattering characteristics based on the canonical model of the tubular cylinder of finite length, more complicated models varying from a tubular cylinder to a missile by adding fins and wings will be constructed and studied. Figures 4.11 and 4.12 show the experimental patterns of a finite tubular cylinder with four fins in two aspect angles, the first, with the fins parallel to the antennas, and the second twisted in 45 degrees. Both Figures can be compared to Figure 3.19 showing the same cylinder without the fins.

This work is the first step to gain the capability of radar target identification. Its contribution lies in the justification that the canonical model is a useful one: the theory is correct so that the theoretical results can provide information about surface current distribution on the cylinder.

Another application of this work could be in the study of the back scattering characteristics of aircraft engines.

Since the signals are modulated, target detection will be easier. And because of the smaller size of the air-intake, the results will be directly appliable to radars now in existance.

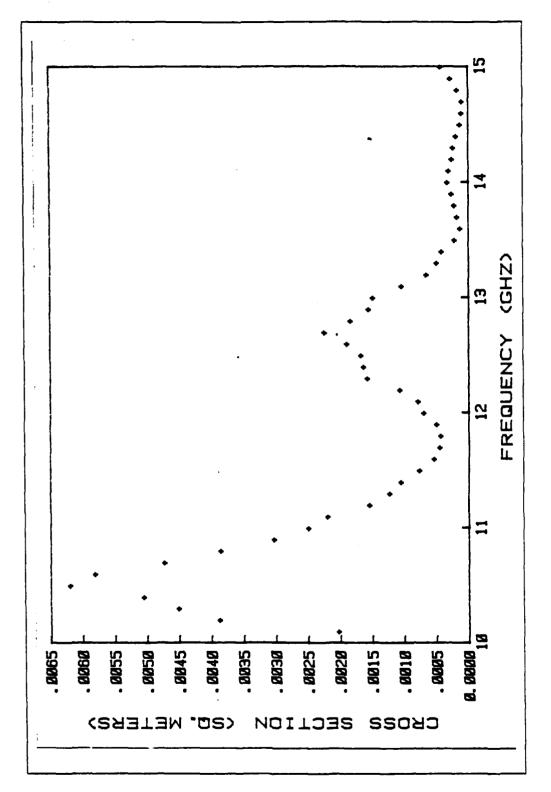
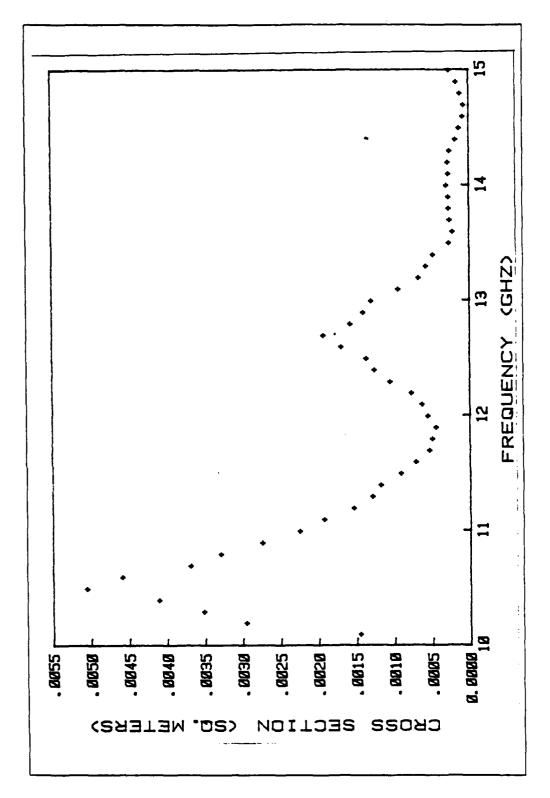


Figure 4.11 Cylinder 15 with Fins, 90 Deg



igure 4.12 Cylinder 15 with Fins,45 Deg

APPENDIX A COMPUTER PROGRAMS

```
10
     ! "SPHERE.DRIVEO"
  20
  30
        COMPUTE BACK SCATTERRED
      ! FAR-FIELD FROM A PERFECTLY
  49
       CONDUCTING SPHERE.
  50
        THE INCIDENT FIELD IS A LINEARLY POLARIZED PLANE W
  60
  70
      AVE WITH ZERO PHASE AT THE CENTER OF THE SPHERE.
      ! THE THEORETICAL VALUES TO
BE COMPUTED ARE THE BACK-SC
      ATTERING CROSS-SECTION AND
     ! THE PHASE OF THE FAR FIELD INTERPOLATED TO THE CENTER
      OF THE SPHERE.
100
110
        THEORETICAL VALUES ARE
        STORED IN FILES OF 800
120
     ! RECORDS; ONE FOR EACH FREQ
UENCY FROM 2.02 GHZ TO 18 GH
Z AT 0.02 GHZ STEPS.
130 !
149
150
        "THEORILDRIVE1" FOR THE
160
        1" DIAMETER SPHERE
        "THEOR3.ORIVE1" FOR THE 3.187" DIAMETER SPHERE
170
180
        "THEOR4.DRIVE1" FOR THE
190
        4.75" DIAMETER SPHERE
200
        "THEORE DRIVE1" FOR THE
219
220
        6" DIAMETER SPHERE
230
240
       FILE H$ STORES THE
     COMPUTED RESULT
260
    H$="THEOR6.DRIVE1"
270
280 A0=6*.0254/2 ! SPHERE
     RADIUS IN METERS
300
310 X9=2*PI ! PARAMETER
330
340 Q1=2 ! STARTING FREQ IN GHZ
350 Q2=18 ! FINAL FREQ IN GHZ
360 Q4=.02 ! FREQ STEP IN GHZ
370 ON ERROR GOTO 390
380 PURGE H≸
390 OFF ERROR
```

```
400 CREATE H$.800.16 ! OPEN A NE
W FILE WITH 800 RECORDS
    ! OF 16 BYTES EACH. EVERY
420 ! RECORD STRORES ONE MAGNITU
     DE AND ONE PHASE DATA.
430
440 ASSIGN# 1 TO H$
450 DIM B8(144),B9(144),D8(144),
     D9(144) ! 144>L0=INT(2*K0*A0
     +3)
460 F0=Q1
470 FOR I=1 TO 800
480 DISP "FREQ LOOP=",I
490 F0=F0+Q4
500 DISP "FREQ (GHZ)=",F0
510 K1=.3/F0 ! WAVELENGTH
520 K0=X9/K1 ! WAVE NUMBER
530 GOSUB 650
540 DISP "E =",E0
550 DISP "P =",P0
560 PRINT# 1,1 ; E0,P0
570 MEXT I
580 ASSIGN# 1 TO *
590 CLEAR
600 DISP "END OF COMPUTATION"
610 END
620
630
    L0=INT(2*K0*A0+3)
640
650 IF L0<145 THEN 670
660 DISP "K0*A0 TOO LARGE FOR
    CURRENT ARRAY DIM"
670 Z=K0*A0
680 GOSUB 910
690 E8=0
700 E9=0
710 FOR N=1 TO L0
720 L=L0-N+1
730 M8=08(L)^2+09(L)^2
740 M9=88(L)^2+89(L)^2
750 A7=(L+.5)/M8/M9
760 A8=A7*(B9(L)*D9(L)-B8(L)*D8(
    L))
770 A9=A7*(B8(L)*D9(L)+B9(L)*D8(
    L))
780 E8=A8-E8
790 E9=A9-E9
800 NEXT N
810 E8=-E8
820 E9=-E3
830 E0=E8^2+E9^2
840 P0=ATN2(E9,E8)
850 E0=E0/K0*2*K1 ! CROSS-SECTIO
   P0=P0-X9*INT(P0/X9)
369
870 IF P0<PI THEN 890
880 P0=P0-X9
```

```
890 P0=-P0
900 RETURN
910
920 IF Z>L0-1 THEN 1230
930 Z2=Z^2/2
940 N2=2*Z2+L0+1
950 D1=2*N2+3
960 D2=D1*(2*N2+5)
970 D3=D2*(2*N2+7)
980 D4=D3*(2*N2+9)
990 F1=1-Z2/D1+Z2^2/(2*D2)-Z2^3/
     (6*D3)
1000 F2=Z*(1/D1-Z2/D2+Z2^2/(2*B3
)-Z2^3/(6*D4))
1010 M=2*Z2
1020 S1=F1
1030 F1=(2*M+1)*F1/Z-F2
1040 F2=S1
1050 IF ABS(F1)<1.E100 THEN 1090
1060 F1=F1*1.E-100
1070 F2=F2*1 E-100
1080 S1=S1*1.E-100
1090 M=M-1
1100 IF M+1>L0 THEN 1020
1110 B8(L0)=F2
1120 B8(L0-1)=F1
1130 NO=L0-2
1140 FOR K=1 TO NO
1150 N=L0-K-1
1160 B8(N)=(2*N+3)*B3(N+1)/Z-B8(
     N+2)
1170 NEXT K
1180 A1=(SIN(Z)/Z-COS(Z))/88(1)
1190 FOR K=1 TO L0
1200 BS(K)=A1*BS(K)
1210 HEXT K
1220 GOTO 1280
1230 B8(1)=SIN(Z)/Z-COS(Z)
1240 B8(2)=(3/Z^2-1)#SIN(Z)-3#C0
     S(Z)/Z
1250 FOR N=3 TO L0
1260 B8(N)=(2*N-1)*B8(N-1)/Z-B8(
     N-2)
1270
    NEXT N
1280 B9(1)=-SIN(Z)-COS(Z)/Z
1290 B9(2)=(1-3/Z^2)*COS(Z)-3*SI
     N(Z)/Z
1300 FOR N=3 TO LO
1310 B9(N)=(2*N-1)*B9(N-1)/Z-B9(
1320 NEXT N
1330 D8(1)=(1-1/2^2)*SIN(2)+COS(
     えシィス
1340 \text{ D9}(1) = (1/2^2-1)*008(2)*SIN(
1350 FOR N=2 TO LO
1360 D8(N)=B8(N-1)-N*B8(N)/Z
```

1370 D9(N)=89(N-1)-N*89(N)/Z 1380 NEXT N 1390 RETURN

```
10
    -! "CALIB DRIVEO"
 20
      CALIBRATION USING A SPHERE OVER L3-U9 GHZ AT F9 GHZ
 30
 40
      STEPS BASED ON THEORETICAL VALUES COMPUTED USING THE P
 50
    ROGRAM "SPHERE DRIVEO"
    ! THE RESULTED SYSTEM TRANS-
     FER FUNCTION IS STORED AS:
 70
       "CALIB3.DRIVE1" (3.187")
 80
 90
      A$ IS THE FILE STORING THE
100
      BACKGROUND DATA.
119
120
130
      C# IS THE FILE STORING THE
      SYSTEM TRANSFER FUNCTION
140
    ! H& IS THE FILE STORING
150
    THEORETICAL DATA OF THE SPHE
    RE
160
      S# DESCRIBES THE SPHERE
179
180
199
200 C$="CALIB3 DRIVE1"
210 H$="THEOR3.BRIVE1"
220 S#="3 187 INCH SPHERE"
230
    A$="BKGRND DRIVE1"
240
250 X9=2*PI ! A PARAMETER
260
270 OPTION BASE 1
280 NO=3 ! NUMBER OF READINGS
    ! TAKEN AND AVERAGED FOR ONE
290
300
      FREQ.
310
    F9= 1 ! FREQ. STEP IN GHZ
320
330
340 M1=51 ! M1=(U9-L9)/F9+2
    ! NUMBER OF FREQ. CHECKED.
360
370 L9=10.1 ! LOWER FREQ IN GHZ
380 U9=15 ! UPPER FREQ.IN GHZ
390
400 DIM A(51,2) ! BACKGROUND
                         DATA
   DIM B(51,2) ! TARGET DATA
DIM G3(51,2) ! THEORY
410
429
    ! CREATE C$,M1,16
430
    ! CREATE A$/M1/16
440
    ! STORE CALIBRATION AND BACK
    GROUND DATA IN A FILE OF MI
    RECORDS
! EACH RECORD CONTAINS ONE
    MAGNITUDE AND ONE PHASE
470 : DATA AT A FREQUENCY
```

```
480
     ! READING THE THEORETICAL DA
 500
 510 ASSIGN# 1 TO H≸
 520 K0=(L9-2-F9*2)*50
 530 FOR I=1 TO M1
 540 K0=K0+50*F9
 550 READ# 1.K0 ; G3(I,1),G3(I,2)
 560 NEXT I
 570 ASSIGN# 1 TO *
580 GOSUB 2230 ! HEADER.
590 DISP "DO YOU WANT TO USE THE
      MOST RECENT BACKGROUND DATA
                            YZN "
 600 INPUT P≇
610 IF P$="H" THEN 690
 620
630
     ! READING BACKGROUND DATA.
649
650 ASSIGN# 4 TO A≸
     RE80# 4 / 8(/)
660
679
    ASSIGN# 4 TO *
    GOTO 940
680
690 CLEAR
 700
    REMOTE 7 ! REMOTE ALL
                       DEVICES
710 CLEAR 7 ! CLEAR ALL DEVICES
720 ! INITIALIZE SIG.GEN TO FIRS
     T FREQ.
730 OUTPUT 719 ,"P",L9,"Z1K0L3M0
     N601"
740 CLEAR
750 DISP "REMOVE TARGET FROM CHA
     MBER, PUSH 'CONT' WHEN READY"
760 LOCAL 7
770 BEEP @ BEEP
780 PAUSE
790 REMOTE 7
800 CLEAR
810 DISP "TAKING BACKGROUND DATA
820 PRINT
830 ! FRINT "BACKGROUND DATA"
840 PRINT
850 OUTPUT 719 ;"P",L9,"Z1K0L3M0
     NE01"
860 WAIT 200 ! WAIT FOR FREQ TO
      STABILIZE
870 GOSUB 1380
889
890
       STORING BACKGROUND DATA
900
910 ASSIGN# 3 TO A$
920 PRINT# 3 , A(,)
930 ASSIGN# 3 TO $
940 CLEAR
```

```
950 LOCAL 7
960 DISP "PUT TARGET INTO CHAMBE
R: PUSH 'CONT' WHEN READY"
 970 DISP "TARGET IS "√S$
 980 BEEP @ BEEF
 990 PAUSE
 1000 REMOTE 7
 1010 CLEAR
 1020 DISP "COMPUTING TARGET DATA
 1030 PRINT
 1040 ! PRINT "TARGET DATA"
 1050 PRINT
      OUTPUT 719 ; "P", L9, "Z1K0L3M
 1060
       0N601"
      WAIT 200
1080 GOSUB 1970
1090 PRINT " "
1100 PRINT "TRANS, FUNCTION",S$
       PRINT
1110
1120
      ! CALCULATE AND STORE TRANS
1130
       FER FUNCTION.
1149
1150 ASSIGN# 2 TO C$
1160 FOR M=1 TO M1
1170 N1=8(M,1)-A(M,1)
      N2=B(M/2)-A(M/2)
%6=G3(M,1)/(N1^2+N2^2)
%7=G3(M,2)-ATN2(N2,N1)
1180
1130
1200
1210 X7=X7-X9*INT(X7/X9)
1220 IF X7>PI THEN X7=X7-X9
1230 PRINT# 2,M ; X6,X7
1240 ! PRINT USING 970 ; M,X6,X7
1250 IMAGE DD,1X,"X6=",SD.DDDE,1
X,"X7=".SD.DDDE
1260 NEXT M
1270 ASSIGN# 2 TO *
1280 CLEAR
1290 DISP "CALIBRATION COMPLETED
       ,DATA STORED IN",C$
1300 BEEP @ BEEP @ BEEP
1310 LOCAL 7
1320
1330
      END
1340
1350
1360
1370
1380
         BACKGROUND DATA COLLECTIO
      N SUBROUTINE
1390
         OUTPUT(L9-F9)TO U9 GHZ AT
1400
F9 GHZ STEPS
1410 J=10*(L9-2*F9) ! FREQUENCY
       STARS AT L9-F9 GHZ
```

```
1420 FOR K=1 TO M1 ! NUMBER OF FREQUENCY STEPS 1430 J=J+10*F9
1440 IMAGE 18.3Z,14A
1450 OUTPUT 719 USING 1440
      ,J,"00ZIKOL3M0N601"
1460 ! TAKE DATA IN FROM 722%720
1470 GOSUB 1640
1430
     ! CALCULATE REAL&IMAGINARY
1490
      FROM AMP. &PHASE
1500
1510 R1=A1*COS(P1)
1520 I1=A1*SIN(P1)
1530 A(K,1)=R1
1540 A(K,2)=I1
1550 ! PRINT USING 2040 ; A(K,1)
      ,A(K,2)
     NEXT K
1560
     OUTPUT 719 ; "P", L9, "Z1K0L3M
1570
     0N601"
1580 RETURN
1590
1600
1610
1620
1630
     ! SUBROUTINE TO ENTER AMPLI
1640
      TUDE AND PHASE DATA FROM DI
      GITAL VOLTMETERS
1650
       PREPARE DIGITAL VOLTMETER
1660
       TO SEND AMPLITUDE DATA
     ! NO READINGS TAKEN AND AVE
RAGED FOR ONE FREQ.
1680
1690
1700 V1=0 ! PARAMETERS FOR THE
1710 F1=0 ! AVERAGING PROCESS.
1720 FOR L=1 TO NO
1730 OUTPUT 720 ; "POFIRITIZIFLOM
     ø"
1740 WAIT 10
1750 ENTER 720 ; V
1760 WAIT 10
1770 OUTPUT 722 ; "F1R7T1M3A0H1"
1780 WAIT 10
1790 ENTER 722 / F
1800 V1=V1+V
1810 F1=F1+F
1820 WAIT 10
1830 NEXT L
1840 V=V1/N0
1850 F=F1/N0
1860 A1=10^F | TRANSFER TO MAG.
       FROM VOLTS.
```

```
1870 P1=100*V ! TRANSFER TO DEG
      FROM VOLTS.
1880 P1=DTR(P1)
1890 ! PRINT USING 1810 ; K.AI,P
1900 IMAGE DD:2X:"A=":MD:DDDE:2X
:"P=":SD:DDDE
1910 ŘETURŃ
1920
1930
1940
1950
1960
1970
      ! TARGET DATA COLLECTION
       SUBROUTINE
1980
        OUTPUT(L9-F9)TO U9 GHZ AT
1990
     F9 GHZ STEPS
J=10*(L9-2*F9) ! FREQUENCY
2000
      STARS AT L9-F9 GHZ
2010 FOR K=1 TO M1 ! NUMBER OF FREQUENCY STEPS.
2020 J=J+10*F9
2030 IMAGE 1A,3Z,14A
2040 OUTPUT 719 USING 2030 ; "P"
      J,"00Z1K0L3M0N6O1"
2050 ! TAKE DATA IN FROM 722%720
      1820 GOSUE 1410
2060 GOSUB 1640
2070 ! CALCULATE REAL&IMAG.FROM
     AMP&FHASE
2080 R1=A1*COS(P1)
2090 I1=A1*SIN(P1)
2100 B(K,1)=R1
2110 B(K,2)=I1
2120 ! PRINT USING 2040 ; B(K,1)
      /B(K/2)
2130 IMAGE 4X, "R=", SD. DDOE, 2X, "I
     ="/SD.DDDE
2140 NEXT K
             719 / "P"/L9/"Z1K0L3M
2150 OUTPUT
     9N601"
2160 RETURN
2170 !
2180
2190
2200
2210
     ! HEADER SUBROUTINE
2220
2230 PRINT "
2240 PRINT
2250 CLEAR
2260 DISP "CALIBRATION STANDARD"
     . S.≇
2270 PRINT "CALIBRATION STANDARD
     ",5$
```

```
16
       "TARGET DRIVEG"
 20
30
       TARGET BACK-SCATTERING
       USING C# DATA & STORE
 40
     RESULTS IN G$
! FREQUENCIES:L9-U9 GHZ AT
      F9 GHZ STEPS
 60
    ! FILE C# STORES THE SYSTEM
 70
      TRANSFER FUNCTION
 80
     ! FILE G# STORES TARGET DATA
 99
      OBTAINED FROM THIS PROGPAM
100
      FILE H& STORES THEORETICAL VALUES FOR PLOTTING OVERLAY
110
120
    ! FILE A≴ STORES BACKGROUND
      DATA
140
150 C$="CALIB3.DRIVE1"
160 G≇="SCR3.DRIVE1"
170 H≸="THEOR3.DRIVE1"
180 A$="BKGRND.DRIVE1"
190
200
      CREATE G$,52,24
       STOPE TARGET DATA IN FILE
216
      OF M1+1 RECORDS FIRST ONE
FOR THE AVERAGE PROCEDURE
AND THE REST CONTAINS THE
220
230
240
    I FREQUENCY MAGNETUDE AND
250
    ! PHASE SHIFT.
260
279
280
    ! CREATE A$,51,16
    ! STORE CALIB. AND BACKGROUN
    DATA IN A FILE OF M1 RECORDS
300 ! EACH RECORD CONTAINS ONE M
    AG AND PHASE AT A FRED
320 OPTION BASE 1
330 NG=2 ! NUMBER OF READINGS
    ! TAKEN AND AVERAGED FOR ONE
340
        FREQUENCY.
350
360 DIM A(51,2) ! BACKGROUND DAT
370 DIM B(51.2) ! TARGET DATA
380 DIM G4(51.2) ! CALIBRATION
390 DIM N(51)3) ! RESULTANT
400 DIM M9(51/3)
410
420 N1=51 !
    I NI=(U9-L9) /F9+2 NUMBER OF
430
440 ! FREG. CHECKED
450 F9= 1 ! FPEQ.STEPS IN GHZ.
```

```
460 U9=15 ! UPPER FREO IN GHZ
470 L9=10.1 ! LOWER FREO. IN GHZ
480 DIM T(800.2) ! STORES THEORE
460 U9=15 ! UPPER FRED
      TICAL DATA
490 X9=2*PI
500
519
        READING TRANSFER FUNCTION
520 i
530 Assīch# i Toʻc$
540 READ# 1 ; G4() 5
550 ASSIGN# 1 TO #
560 ! MAT PRINT USING 330 ; G4
570 IMAGE 2%,30.40
580 1
590 REMOTE 7 ! REMOTE ALL
                          DEVICES
600 CLEAR 7 ! CLEAR ALL DEVICES
610 OUTPUT 719 : "P1Z1K0L3M0N601"
! INITIAL SETUP OF 719
620 CLEAR
630 DISP "DO YOU WANT TO USE THE
      MOST RECENT BACKGROUND DATA
649 INPUT P#
650
560 IF P#="N" THEN 740
630
     ! READING BACKGROUND DATA
590
700 ASSIGN# 4 TO A#
710 PEAD# 4 ; A(,)
720 ASSIGN# 4 TO *
730 GOTO 900
740 DISP "REMOVE TARGET FROM
     CHAMBER, PUSH 'CONT' WHEN REA
     DY"
750 LOCAL 7
760 BEEP @ BEEP
770 PAUSE
780 DISP "TAKING BACKGROUND DATA
790 REMOTE 7
800 OUTPUT 719 /"P"/L3,"Z1KGL3MG
     NGO1" ! INITIAL SETUP OF 719
810 WAIT 100
820 GOSUB 2560
830
849
     ' STORING BACKGROUND DATA.
850
860 ASSIGN# 5 TO A$
870 PRINT# 5 ; AC, 7
880 ASSIGN# 5 TO *
890 CLEAR
    DISP "PUT TARGET INTO CHAMBE
F, PUSH 'CONT' WHEN READY"
900
219
    LOCAL 7
920 BEEP @ BEEP
930 PAUSE
```

```
940 REMOTE 7
950 OUTPUT 719 ;"P",L9,"Z1K@L3M0
N601" ! INITIAL SETUP OF 719
 960 WAIT 500
 970 GOSUB 3410
 980 CLEAR
990 DISP "COMPUTING TARGET DATA"
 1000 GOSUB 3136
 1010
 1020 ! COMPUTING TARGET DATA
 1030 ! WITHOUT BACKGROUND AND TH
 1040 T FREQ. FOR EACH RECORD.
1050 F0=L9-2*F9
 1060 FOR M=1 TO N1
1070 F0=F0+F3
 1980 N(M,1)=F0
 1090 X7=8(M,1)-A(M,1)
 1100 X8=8(M,2)-A(M,2)
 1110 X6=(X7^2+X8^2)*G4(M,1)
1120 N(M,2)=X6
1130 X8=ATN2(X8,X7)+G4(M,2)
 1140 X8=X8-X9*INT(X8/X9)
 1150 IF X8>PI THEN X8=X8-X9
 1160 \text{ N(M,3)=x8}
 1170 NEXT M
 1180 DISP "PRINT DATA? Y/N"
 1190 BEEP @ BEEP
1200 INPUT P#
1210 IF P#="H" THEN 1250
1220 PRINT "
                       FREQ
                                     CRSEC
            PHASE"
1230 MAT PRINT USING 1240 ; N
1240 IMAGE 28,30,40
1250 CLEAR
1260 LOCAL
1270 DISP "PLOT MAGNITUDE FOR
              THIS MEASURMENT? YOH"
1280 INPUT PE
1290 IF P#="H" THEN 1330
1300 DISP "SELECT PEN PUSH 'CONT' WHEN READY"
1310 PAUSE
1320 GOSUB 3570
1330 CLEAR
1340 DISP "PLOT PHASE FOR THIS
             MEASURMENT ? YZH"
1350 BEEP @ BEEP
1350 BEEP & DLL.
1360 INPUT P$
1370 IF P$="N" THEN 1410
1380 DISP "SELECT PEN. PUSH
__'CONT' WHEN READY"
1400 GOSUB
              4619
1410 CLEAP
```

```
1420 DISP "DO YOU WANT TO MAKE
           AVERAGE WITH PREVIUSE
      DATA?"
1430 DISP "?Y/N"
     BEEP @ BEEP
1440
1450 INPUT P#
1460 IF P#="7" THEN 1750
1470 DISP "DO YOU WANT TO STORE
           DATA ? Y/N
1480 INPUT P#
1490 IF P$="N" THEN 2280
1500 M0=1
1510 DISP "DO YOU WANT TO STORE DATA IN FILE"
1520 DISP G#
1530 DISP "? Y/N"
1540 INPUT P$
1550 IF P$="Y" THEN 1620
1560 DISP "ENTER NAME OF THE DAT
A_FILE TO BE USED FOR STORE
     GE"
1570 INPUT G#
1580 DISP "IS THIS AN OLD FILE
TO BE UPDATED ? Y/N "
1590 INPUT P$
1600 IF P#=""" THEN 1620
1610 CREATE G$.53,24
1620 DISP "ENTER LENGTH OF TARGE
1630 BEEP @ BEEP
1640 INPUT M1
1650 DISP "ENTER DIAMETER OF TAR
     GET"
1660 BEEP @ BEEP
1670 INPUT M2
1680
       STORE MEASURED DATA.
1690
1740 GOTO 2280
1750 DISP "DOES THE DATA STORED
            IN FILE"
1760 DISP G#
1770 DISP "? Y/N"
1780 INPUT P#
1790 IF P#="Y" THEN 1970
1300 DISP "ENTER NAME OF DATA
          FILE TO BE USED FOR THE
           AVERAGE"
1810
1820
     INPUT G$
1830
1340
        READ OLD DATA
1850
        AND MAKES WIGHTED AVERAGE
     ! WITH HEW DATA.
1360
```

```
1870 ASSIGN# 6 TO G$
1880 READ# 6 : M0.M1.M2.M9(,)
1890 ASSIGN# 6 TO *
1900 FOR K=1 TO NI
1910 M9(K,2)=M9(K,2)*M0+N(K,2)
1920 M9(K,2)=M9(K,2)/(M0+1)
1930 M9(K,3)=M9(K,3)*M0+N(K,3)
1940 M9(K,3)=M9(K,3)/(M0+1)
1950 N(K,2)=M9(K,2)
1960 N(K,3)=M9(K,3)
1970 NEXT K
1980 M0=M0+1
1990
2000 ! STORE NEW AVERAGE
2010 ASSIGN# 7 TO G*
2020 PRINT# 7 ; M0.M1.M2.N(;)
2030 ASSIGN# 7 TO *
      PRINT "DATA IS AVERAGE OF".
2040
M0."MEASURMENTS"
2050 DISP "PRINT DATA? Y/N"
2060 BEEP @ BEEF
2070 INPUT P#
2080 IF P$="N" THEN 2120
2090 PRINT "
                        FREQ
                                       CRSEC
            PHASE"
2100 MAT PRINT USING 1240 ; N
2110 IMAGE 2X/3D.4D
2120 DISP "PLOT MAGNITUDE?
2130 BEEP @ BEEP
2140 INPUT P$
2150 IF P$="N" THEN 2190
2160 DISP "SELECT PEN FOR MAGNIT
UDE PLOT PUSH 'CONT' WHEN
       READY.
2170 PAUSE
2180 GOSUB 3570
2190 CLEAR
2200 DISP "PLOT PHASE? Y/N"
2210 BEEP @ BEEP
2220 INPUT P#
2230 IF P#="N" THEN 2270
2240 DISP "SELECT PEN AND CHANGE
PAPER FOR PHASE PLOT PUS
       H 'CONT' WHEN READY.
2250 PAUSE
2260 GOSUB 4610
2270 CLEAR
2280 DISP "DO YOU WANT TO"
2290 DISP "OBTAIN DATA"
2300 DISP "FOR A NEW TARGET?"
      DISP "
2310
       DISP "ENTER Y/N"
2320
2330
      INPUT PE
2340
      IF P#="N" THEN 2470
2350 DISP "DO YOU WANT TO USE
        THE SAME FILE"
2360 BISP G#
```

```
2370 DISP "TO STORE NEW DATA? Y/
      H"
2380 INPUT P#
2390 IF P$="Y" THEN 2460
2400 DISP "ENTER NEW FILE NAME T
      O STORE TARGET DATA"
2410 INPUT G$
3420 DISP "IS THIS AN OLD FILE
      TO BE UPDATED? Y/Nº
2430 INPUT P$
2440 IF P$="Y" THEN 2460
2450 CREATE G$,52,24
2460 GOTO 590
2470 CLEAR
2480 DISP "END OF PROGRAM"
2490 BEEP @ BEEP @ BEEF
2500 END
2510
2520
2530
2540
2550
2560
        BACKGROUND DATA COLLECTIO
      N SUBROUTINE
2570
      ! OUTPUT(L9-F9)TO U9 GHZ
      J=10*(L9-2*F9) ! FREQUENCY
2580
      STARTS AT L9-F9 GHZ TO BE
INCREASED AT F9 GHZ STEPS
2590 FOR K=1 TO N1 ! NUMBER OF F
      REQUENCY STEPS
2600 J=J+10*F9
2610 IMAGE 1A,3Z,14A
2620 OUTPUT 719 USING 2610 ;
      )J,"00Z1K0L3M0N6O1"
2630 ! 50 MSEC WAIT FOR FREQUENC
Y TO STABILIZE
2640 WAIT 50
2650 ! TAKE DATA IN FROM 722 AND
       720
2660 GOSUB 2820
2670 ! REAL AND IMAGINARY PARTS
2680 ! FROM AMP. AND PHASE.
2690 R1=A1*COS(P1)
2700 I1=A1*SIN(P1)
2710 A(K,1)=R1
2720 A(K,2)=I1
2730 ! PRINT "I1=",A(K,2)
2740 ! PRINT "R1=",A(K,1)
2750 NEXT K
2760 OUTPUT 719 ; "P", L9, "Z1K0L3M
0N601" ! INITIAL SETUP OF 7
2770 RETURN
2780
2790
```

```
2899 !
2810 !
2820 ! SUPROUTINE TO ENTER AMPLI
      TUDE AND PHASE DATA FROM DI
      GITAL VOLTMETER
2830
     ! PREPARE DIGITAL VOLTMETER
2840
       TO SEND AMPLITUDE DATA
2850 ! NO READINGS TAKEN AND AVE RAGED FOR ONE FREQUENCY.
2860 V1=0 ! PARAMETERS FOR AVERA
                  GING PROCESS.
2870 W1=0
2880 FOR L=1 TO NO
2890 OUTPUT 720 /"F1R1T1Z1FL0M0"
2900 WAIT 10
2910 ENTER 720 ; VO
2920 WAIT 10
2930 OUTPUT 722 ;"F1R7T1M3A0H1"
2940 WAIT 10
2950 ENTER 722 ; WO
2960 V1=V1+V0
2970 W1=W1+W0
2980 WAIT 10
2990 NEXT L
3000 V0=V1/N0
3010 W0=W1/N0
3020 ! TRANSFERS FROM VOLTS TO
           AMPL.
3030 A1=10^W0
3040 ! TRANSFERS TO DEG. FROM VO
            LTS
3050 Pi=100*V0
3060 ! PRINT "A1=",A1
3070 P1=DTR(P1)
3080 ! PRINT "P1=".P1
3090 RETURN
3100
3110
3120
3130
     - 1
        DATA COLLECTION SUBROUTIN
      ! OUTPUT(L9-F9)TO U9 GHZ AT
      F9 GHZ STEPS
U=10*(L9-2*F9) !
3150
                         INITIAL FR
      EQUENCY AT L9-F9 GHZ
3160 FOR K=1 TO N1 ! FREQUENCY S
      TEPS
3170
      J=J+10*F9 ! F9
                         GHZ INCREME
      NTS
3180 IMAGE 18,32,148
3190 OUTPUT 719 USING 3180 ;
,J,"00Z1K0L3M0N601"
      ! 50 MSEC WAIT FOR FREQUENC
      Ý TO STABILIZE
3210 WAIT 50
```

```
3220 ! TAKE DATA IN FROM 7221720
3230 GOSUB 2820
3240 !
3250 ! REALGIMAG FROM AMP.&PHASE
3260 R1=A1*COS(P1)
3270 I1=A1#SIN(P1)
3280 B(K,1)=R1
3290 B(K,2)=I1
3300 ! PRINT "R1=".B(K,1)
3310 ! PRINT "I1=".B(K,2)
3320 NEXT K
3330 OUTPUT 719 ; "P", L9, "Z1K0L3M
0N601" ! INITIAL SETUP OF 7
3340 RETURN
3350
3360
3370
      ! HEADER SUBROUTINE
3380
3390
3400 D#="MONTH/DATE/YEAR"
3410 PRINT
3420 PRINT
3430 CLEAR
3440 DISP "ENTER TODAY'S DATE -
      MONTH, DATE, YEAR"
3450 INPUT D$
3460 DISP "ENTER TGT DESCRIPTION
3470 INPUT T≉
3480 PRINT 0#
            "TARGET IS ",T$
3490 PRINT
             "*************
3500 PRINT
3510 PRINT "*************
3520 PRINT
3530 CLEAR
3540 RETURN
3550
3560
3570
        MAGNITUDE PLOTTING
3580
        SUBROUTINE
3590
3600 PLOTTER IS 705
3610 LOCATE 32,122,20,85
3620 FRAME
        SEARCH FOR MAX
3630
                          & MIN.
3640
        S0=N(2,2)
3650
        31=90
3660
     ! FOR M=3 TO N1
     ! IF $00N(M,2) THEN $0=N(M,
3670
3630
     ! IF SIKN(M)2) THEN SI=N(M)
      2)
3690 ! NEXT N
           "ENTER LOWER VALUE FOR
3700 DISP
           MAGNITUDE PLOTTING"
```

```
3710 BEEP @ BEEP
 3720 INPUT SO
3730 DISP "ENTER UPPER VALUE FOR
             MAGNITUDE PLOTTING"
 3740 INPUT SĪ
 3750 L1=INT(L9)
 3760 U1=CEIL(U9)
 377Ø
 3780
      ! CALCULATE SCALE STEPS
 3790 ! FOR MAGNETUDE.
 3800 S3=LGT(S1)
 3810 S4=INT(S3)-1
 3820 85=83-84
3830 85=INT(10^85)+1
 3840 84=10^84
 3850 L0=1NT(80/84)
 3860 IF S5-L0<=14 THEN 3940
 3870 IF S5-L0>=50 THEN 3910
 3880 55= 5#85
 3890
      34=$4*2
 3900 GOTO 3850
3910 85= 2*85
3920 84=5*84
3930 GOTO 3850
3940 L0=84*L0
3950 00=95*84
3960 D0=00-L0
3970 SCALE L1,U1,L0,U0
3980 FXD 0,4
3990 LAXES -1,84,L1,L0
4000 MOVE L1,0
4010 FOR K=2 TO N1
4020 N9=N(K,1)
4030 R9=N(K, 2)
4040 GOSUB 4510
4050 NEXT K
4060 M5=(U1+L1)/2
4070 MOVE M5.L0- 09*D0
4080 LORG 5 @ CSIZE 3.1.0
4090 ! LABEL "FREQUENCY (GHZ)"
4100 MOVE L1-1,.5*(L0+U0)
4110 LDIR PI/2
4120 ! LABEL "CROSS SECTION (SQ. METERS)"
4130 MOVE M5/U0+.09*D0
4140 LDIR 0
4150 CSIZE 3,1,0
4160 LABEL T#
4170 MOWE M5.U0+.03*D0
4180 LABEL D≇
4190 PENUP
4200 DISP "OVERLAY THEORETICAL C
     URME? Y'N"
4210
4220
     INPUT P$
4230 IF P$="N" THEN 4496
```

```
4240 DISP "IS THE THEORETICAL DATA STORED IN THE FILE", H$ 4250 DISP "? YZN"
      INPUT P#
4260
      IF P$="Y" THEN 4366
4270
4280 GISP "ENTER NAME OF THE DATA FILE TO BE PLOTTED."
4290 INPUT H$
4300 BEEP @ BEEP
4310 DISP "CHANGE PEN IF DESIRED PUSH 'CONT' WHEN READY."
4320 PAUŠĒ
4330 ASSIG
      ASSIGN# 3 TO H$
4340 J1=(L9-2)*50
      J2=(U9-2)*50
4350
4360 FOR J=J1 TO J2
4370 READ# 3.J ; T(J,1),T(J,2)
4380 NEXT J
     ASSIGN# 3 TO *
4390
4400
      F0=L9
      R9=T(J1,1)
4410
4420 MOVE F0, R9
4430 FOR I=J1+1 TO J2
4440 F0=F0+.02
4450 R9=T(I,1)
4460 DRAW F0,R9
4470 NEXT I
4480 PENUP
4490 RETURN
4500
4510
        PLOT CROSS
4520 MOVE M9.R9
4530 CSIZE 2, 5,0
4540 LABEL "+"
      ! IMOVE .00025..00025
! IDRAN -.0005.0
4550
4560
4570
      RETURN
4580
4590
4600
4610
      ! PHASE PLOTTING SUBROUTINE
4620
4630 PLOTTER IS 705
4640 LOCATE 32,122,20,85
4650 FRAME
      84=.25
4660
4670 00=1
4680 L0=-1
4690 D0=U0-L0
4700 SCALE L1,U1.L0,U0
     FXD 0.3
4710
4720
4730
      LAXES -1/84/L1/L0
4730 NOVE L1,0
4740 FOR K=2 TO N1
4750 M9=N(K,1)
4760 R9=N(K,3) PI
4770 GOSUB 5110
```

```
4780 NEXT K
4790 MOVE M5.L0-.09*D0
4800 LORG 5 @ CSIZE 3:1:0
       ! LABEL "FREQUENCY (GHZ)"
4810
4820 MOVE L1-1. 5*(L0+U0)
4830 LDIR PI/2
4840 LABEL "PHASE (PI)"
4850 MOVE M5,00+.09*00
4860 LDIR 0
4870 CSIZE 3,1,0
4880 LABEL T#
4890 MOVE M5, U0+.03*00
4900 LABEL D#
4910 PENUP
4920 DISP "OVERLAY THEORETICAL C
URVE? Y/N"
4930 INPUT P#
4940
       IF P$="M" THEN 5070
4950 BEEP @ BEEP
4960 DISP "CHANGE PEN IF DESIRED
. PUSH 'CONT' WHEN READY."
4970 PAUSE
4980 F0=L9
4990 R9=T(J1/2)/PI
5000 MOVE F0.R9
5010 FOR I=J1+1 TO J2
5020 F0=F0+.02
5030 R9=T(I,2)/PI
5040 DRAW F0,R9
5050 NEXT I
5060 PENUP
5070
       RETURN
5080
5090
5100
          PLOT DOT
5110
5120 MOVE MOJRS
5130 CSIZE 2, 5.0
5140 LABEL "*"
      ! IMOVE .00025,.00025
! IDRAW -.0005,0
5150
5160
5170 RETURN
```

$\frac{\texttt{APPENDIX} \ \, \texttt{B}}{\texttt{CHRACTERISTIC} \ \, \texttt{OF} \ \, \texttt{THE} \ \, \texttt{ABSORBER} \ \, \texttt{OF} \ \, \texttt{THE} \ \, \texttt{ANECHOIC} \ \, \texttt{CHAMBER}}$

MECHANICAL SPECIFICATIONS:

Pyramids	Per Absorber	64
Pyramid Base	Size (In.)	3×3
(In.)	Pyrainid	2
Absorber Height (In.)	Base	1-1/3
Absor	Overall	8-1/2
Absorber	Size (In.)	24 × 24

MAXIMUM REFLECTION AT NORMAL INCIDENCE:

3 ¥	×	U	\$	1	200	300	200	120
 Band	Band	Band	Band	Band	MHz	MHz	MHz	MHz
50 db	db 02	45 db	40 db	30 db				

Specifications of the Absorber Material of the Floor

MECHANICAL SPECIFICATIONS:

Absorber	Abso	Absorber Height (In.)	ıt (ln.)	Pyrainid Base	Pyramids
Size (In.)	Overall	Bose	Pyramid	Size (In.)	Per Absorber
24 × 24	18-1/4	2-1/4	16	9×9	91

MAXIMUM REFLECTION AT NORMAL INCIDITICE:

120	MHZ	
200	Mriz	
300	MHZ	
200	MFIZ	30 db
	pana	40 db
S.	pana	45 db
ں ٔ	Band	50 db
×	pulpa	50 db
Y.	pava	50 db

Specifications of the Absorber Material of the Walls

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